

Dark matter search with coherent atoms

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Outline

- ❑ Introduction
- ❑ Macro-coherent amplification
- ❑ Cesium E2 transition measurement
- ❑ Future Prospect and summary

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❑ **Introduction**

❑ **Macro-coherent amplification**

❑ **Cesium E2 transition measurement**

❑ **Future Prospect and summary**

Introduction

- ❑ Nature of Dark Matter (DM) is still a big mystery.

- ❑ Dark matter candidates

WIMPs, sterile neutrino...

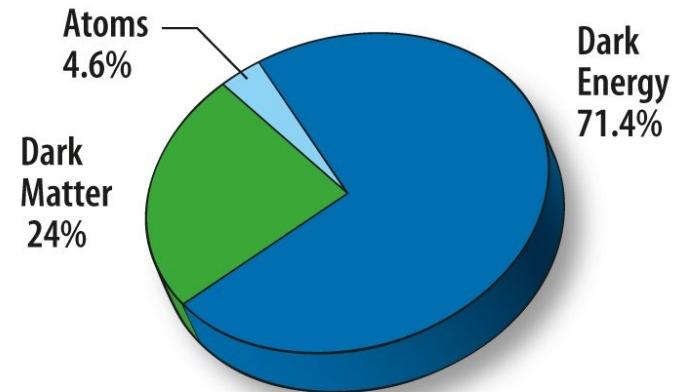
However, null results are found.

- ❑ Axion

$$\mathcal{L}_a = \frac{1}{2}(\partial_\mu a)^2 - \frac{1}{2}m_a^2 a^2 - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{4}g_{a\gamma\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu} \quad g_{a\gamma\gamma} = g\frac{\alpha}{\pi f_a}$$

- ❑ Dark photon

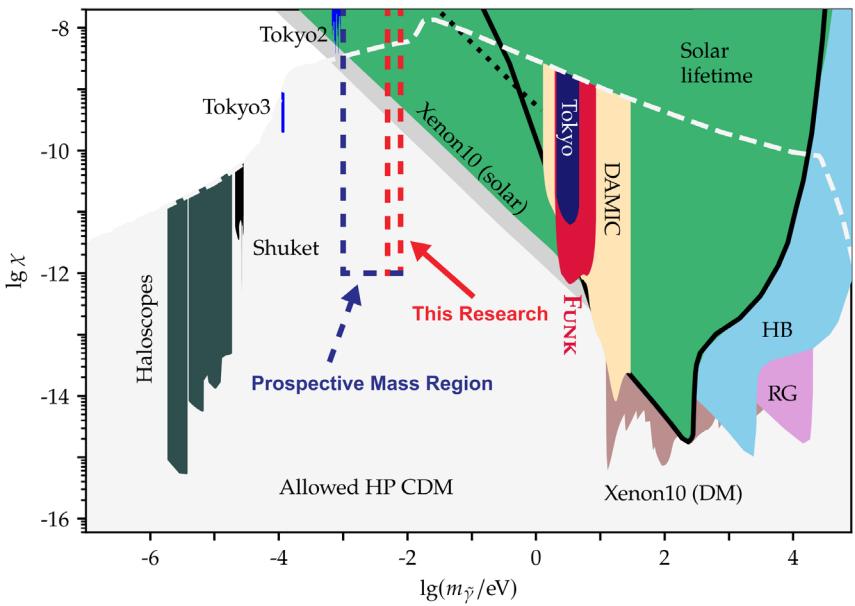
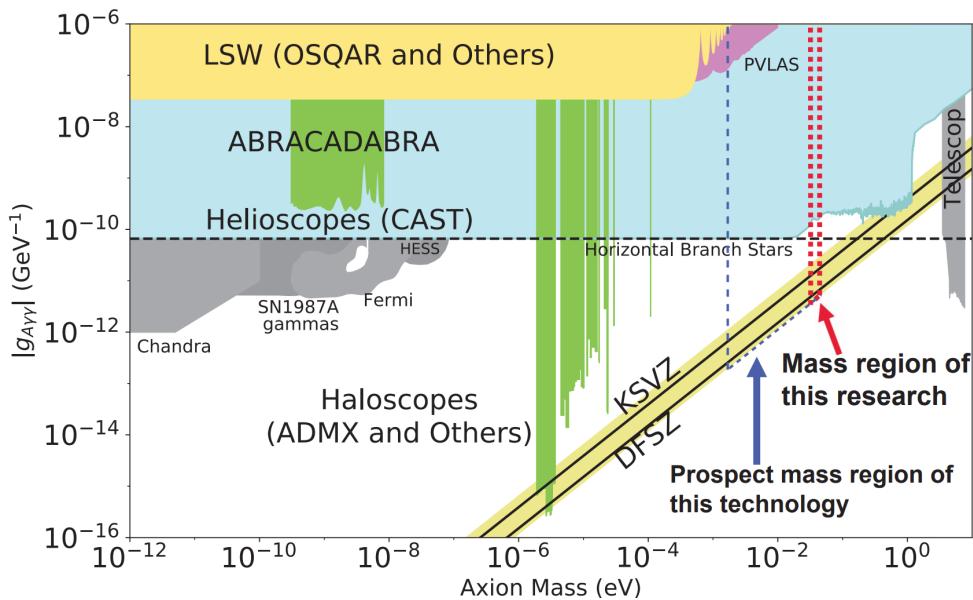
$$\mathcal{L}_{DP} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} - \frac{\chi}{2}F_{\mu\nu}X^{\mu\nu} + \frac{m_X^2}{2}X_{\mu\nu}X^{\mu\nu} + j_\mu A^\mu$$



Axion/Dark photon direct detection with atomic transition

Introduction

- Searching for DM with mass \sim meV



- Amplification mechanism is required

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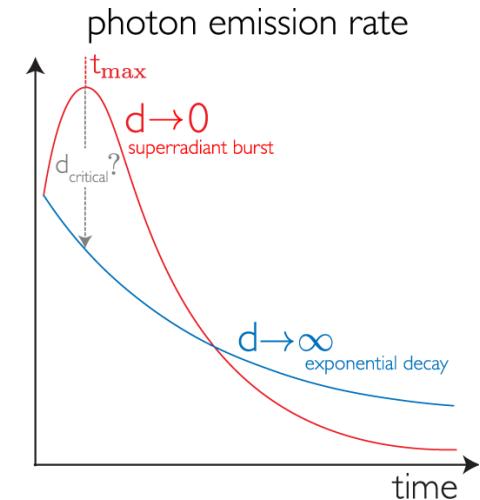
Macro-coherent amplification

□ Super-Radiance [R. H. Dicke Phys. Rev. 93, 99–110 (1954)]

De-excitation via single-photon emission. $|i\rangle \rightarrow |f\rangle + \gamma$

$$R \propto \left| \sum_{m=1}^{N_T} \text{Exp}\left(i\vec{k}_\gamma \vec{x}_m\right) M(\vec{x}_m) \right|^2 \propto N_T^2 \quad [\because M(\vec{x}_m) = M(0), |\text{target size}| \ll \lambda]$$

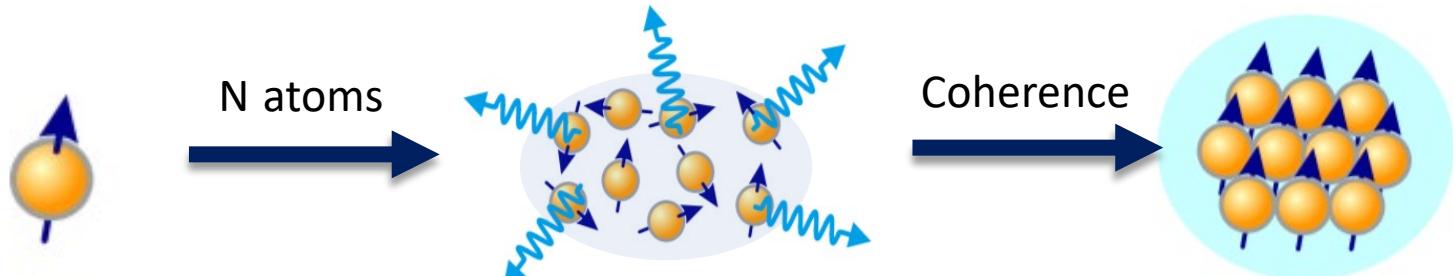
$$\text{Exp}[0] = 1$$



□ Macro-coherent amplification

Multi-photon involved transition. $|g\rangle + \gamma_1 \rightarrow |i\rangle \rightarrow |f\rangle + \gamma_2 \dots$

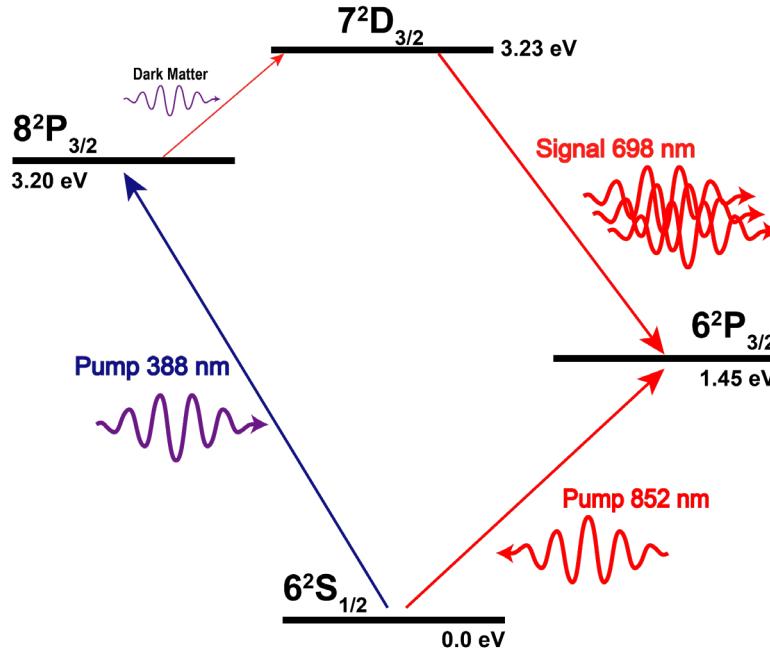
$$R \propto \left| \sum_{m=1}^{N_T} \text{Exp}\left(i\Delta\vec{k}x_m\right) M(\vec{x}_m) \right|^2 \propto N_T^2 \quad [\because M(\vec{x}_m) = M(0), \Delta\vec{k} = -\vec{k}_{\gamma 1} + \vec{k}_{\gamma 2} + \dots = 0]$$



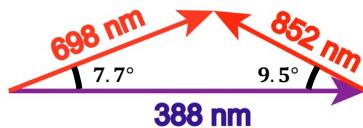
[Miyamoto et al. Prog. Theor. Exp. Phys. 081C01 (2015)]

Dark matter detection scheme

- DM absorption by cesium atom [Sasao et al. Eur. Phys. J. C (2018) 78:949]



- To satisfy $\vec{\Delta k} = 0$

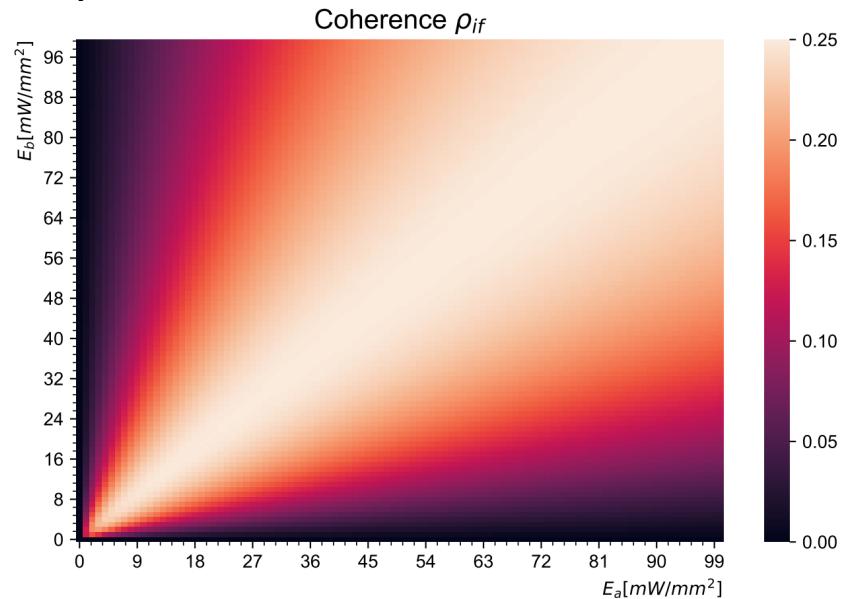
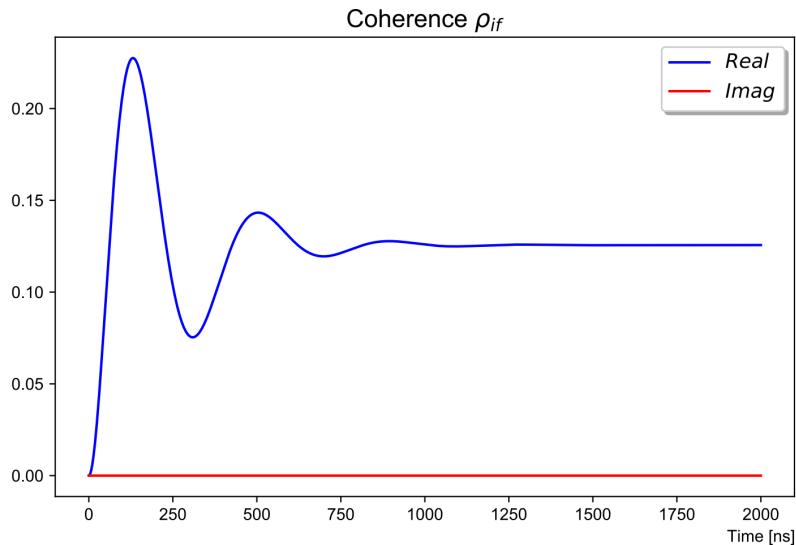


Coherence generation

- The Von Neumann equation

$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}] + G\hat{\rho}G^\dagger - \frac{1}{2} \{ G^\dagger G, \hat{\rho} \}$$

- Inject two pump cw lasers: system reach steady state



- Laser power requirements

$$P_a = \frac{10^{10} \times 3.198 \text{ eV}}{274 \text{ ns}} \sim 20 \text{ mW} \quad P_b = \frac{10^{10} \times 1.5 \text{ eV}}{35 \text{ ns}} \sim 70 \text{ mW}$$

DM signal rate

- ☐ Fermi's Golden Rule for N atoms system

$$d\Gamma = \frac{2\pi}{\hbar} \delta(\hbar(\omega_{fi} + \omega_s - \omega_d)) |M_0 F_{QC}|^2 \frac{V}{(2\pi)^3} d^3 \vec{k}_s$$

- ☐ For finite target (1 mm^3), F_{QC} is described by shape function

$$F_{QC} = N_T |\rho_{fi}| \sqrt{A(\Delta \vec{k})}$$

- ☐ And assuming the following parameters

$$N_T |\rho_{fi}| = 10^{10}, \quad \int A d\Omega = 10^{-7}, \quad T = 4 \text{ K},$$

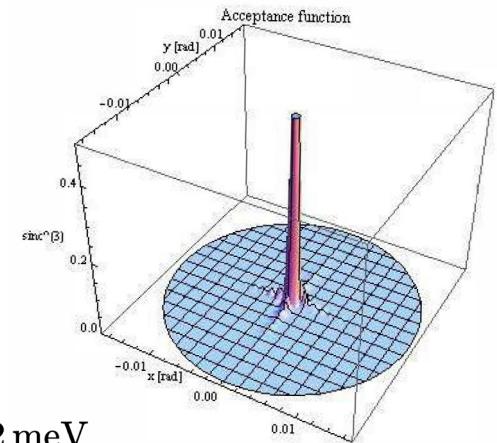
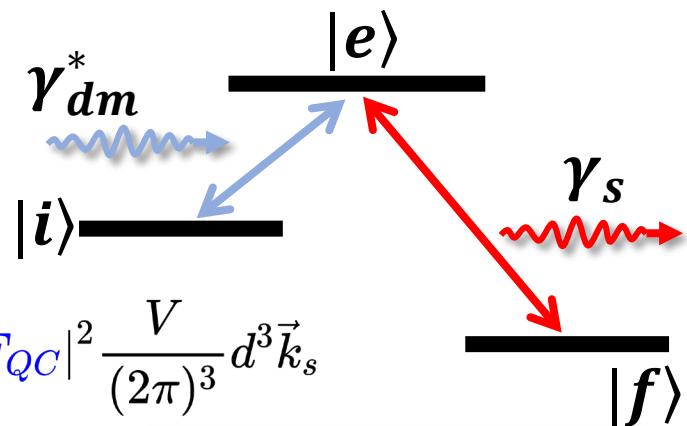
$$\rho_{DM} = 0.3 \text{ GeV/cm}^3, \quad \vec{B}_{ext} = 1 \text{ T}, \quad m_{dm} = 32 \text{ meV}$$

- ☐ For dark photon, $\chi = 10^{-12}$

$$E_d = 3.3 \times 10^{-9} \text{ V/m} \quad \Gamma_{\text{Dark}} \simeq 12.6 \text{ Hz}$$

- ☐ For Axion, $g_{a\gamma\gamma} = 10^{-12} \text{ GeV}^{-1}$, need to increase N_T by 10^4

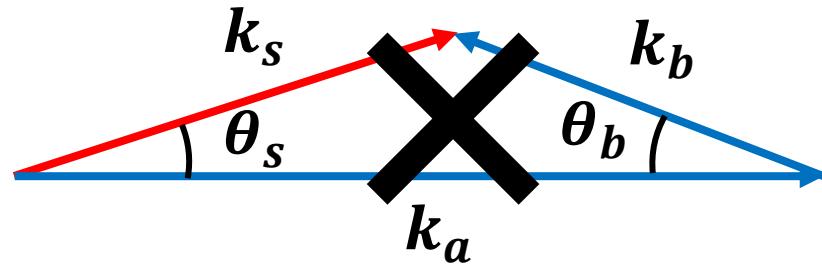
$$E_d \simeq 10^{-13} \text{ V/m} \quad \Gamma_{\text{Axion}} \simeq 10 \text{ Hz}$$



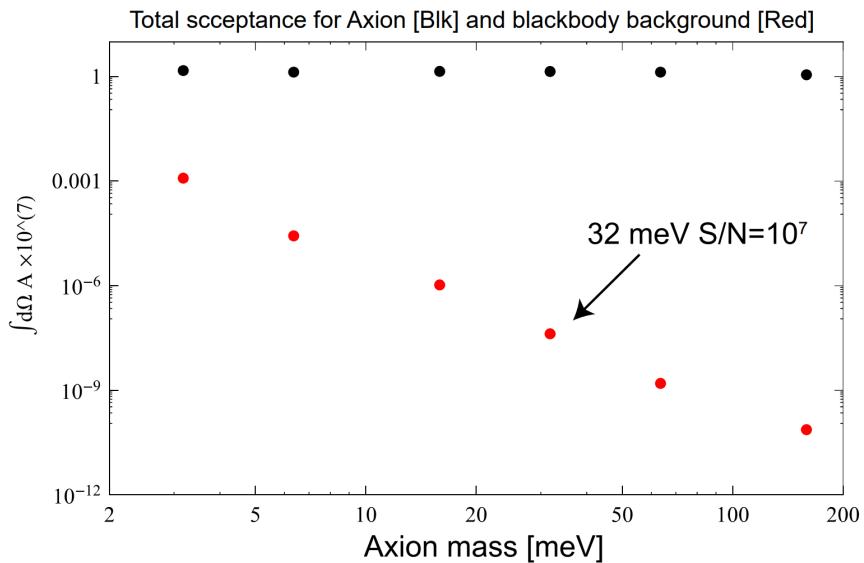
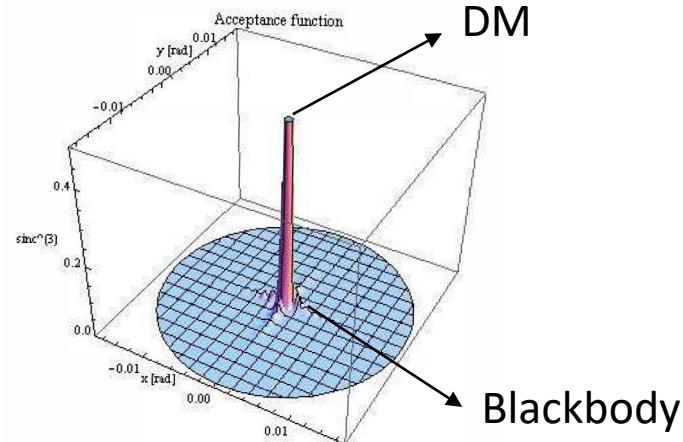
Blackbody background

- Blackbody radiation field strength ($T = 10 \text{ K}$, $m_a \approx 32 \text{ meV}$)

$$E_{bb} \simeq 1.4 \times 10^{-9} \text{ V/m}$$

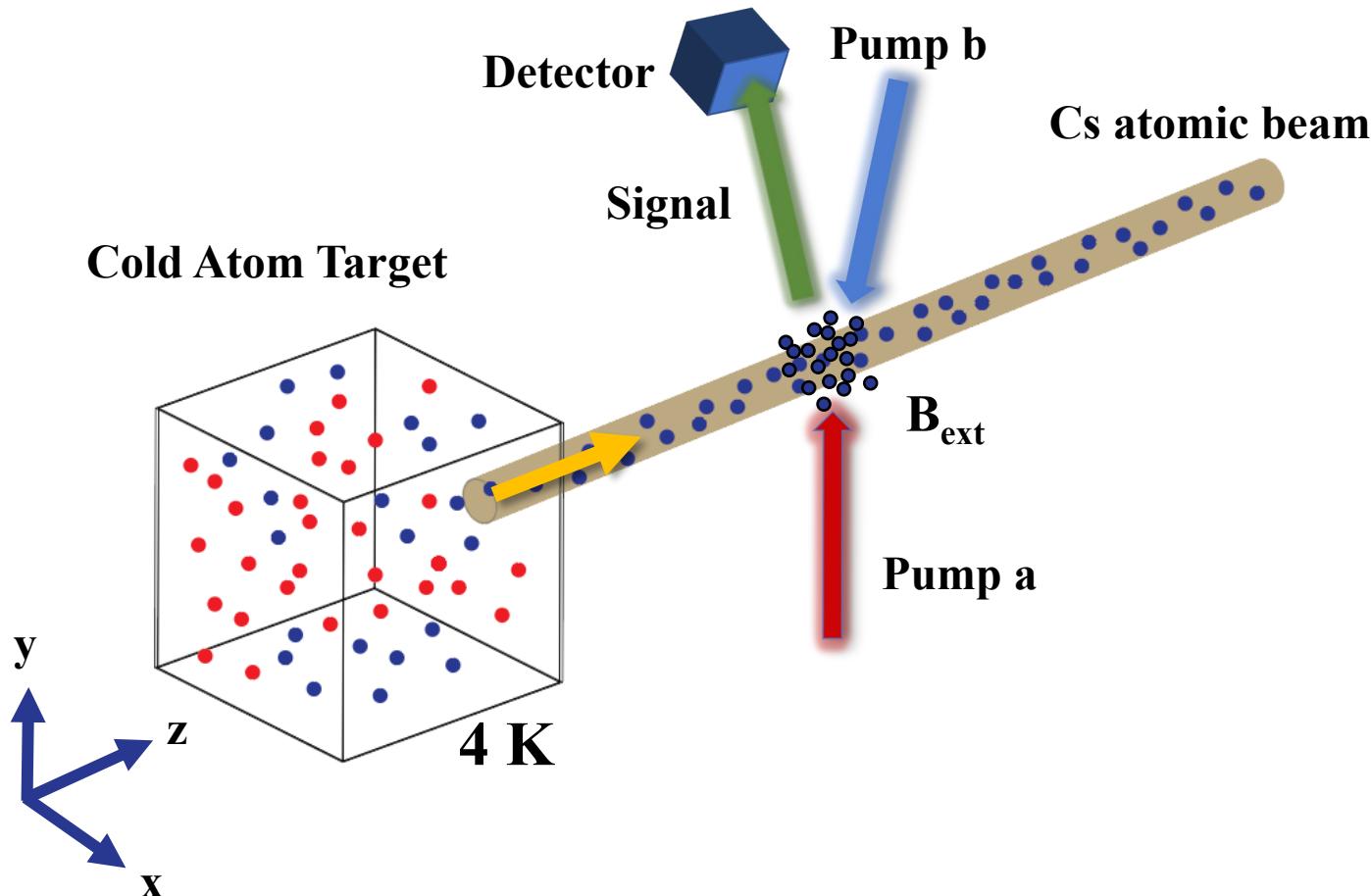


- Blackbody radiation Δk is larger, F_{QC} is small



- Cool down to $T < 20 \text{ K}$

Experiment setup



- Signal rate: ~ 10 Hz
- Temperature: ~ 4 K
- Detector dark count: < 1 Hz
- Mass scan range: $31.8^{+0.0}_{-0.3}$ meV
- Initial coherence: $|\rho_{fi}| > 0.2$
- Laser line width: < 100 kHz

Outline

❑ Introduction

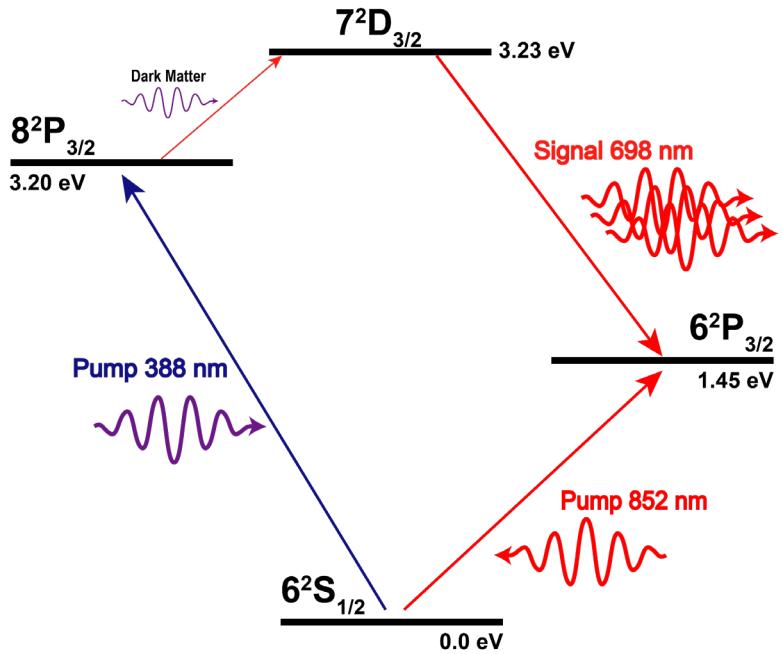
❑ Macro-coherent amplification

❑ Cesium E2 transition measurement

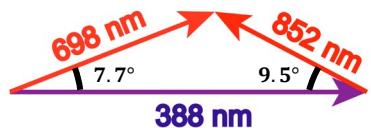
❑ Future Prospect and summary

Coherence evaluation

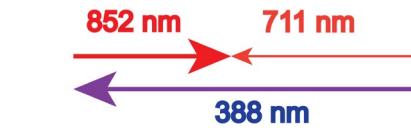
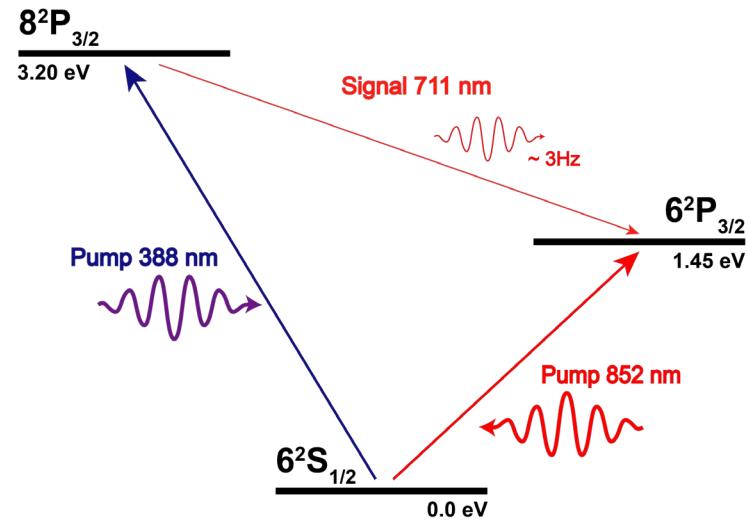
- Using cesium E2 transition



- To satisfy $\Delta \vec{k} = 0$

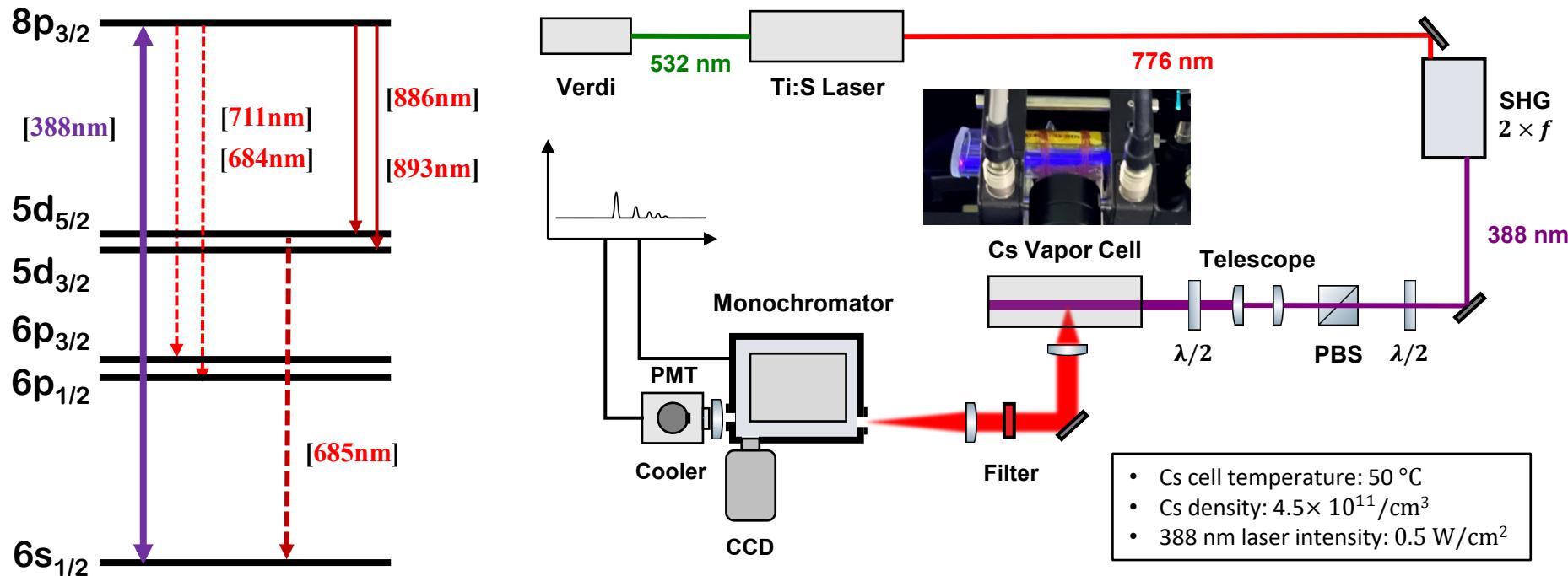


DM detection



Evaluate the coherence

Spectroscopy Setup

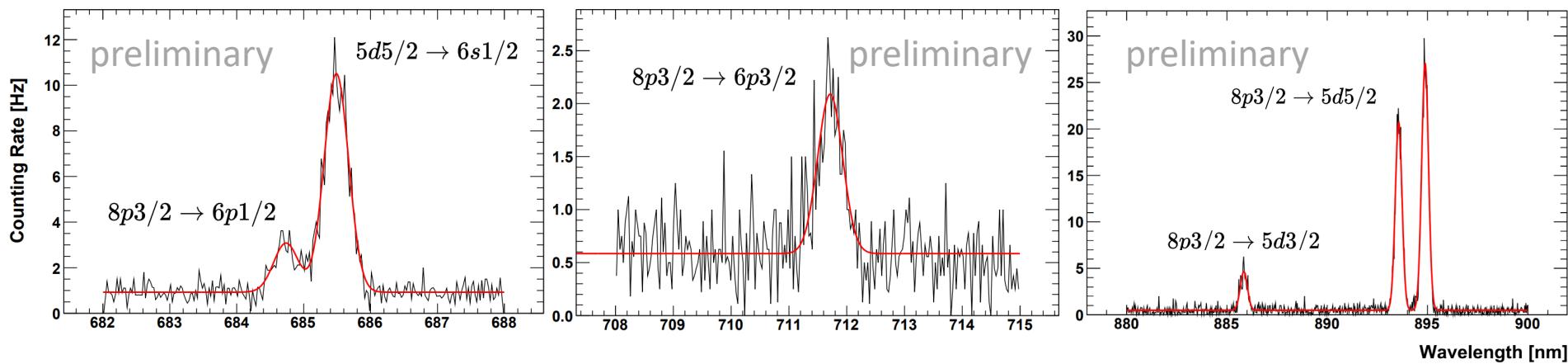


- Compare the emission rate from same excited state
- Photomultiplier Tube (PMT Hamamatsu R13456P) is cooled down to -30 °C to reduce the dark count rate (~1 Hz)

- Cs cell temperature: 50 °C
- Cs density: $4.5 \times 10^{11} / \text{cm}^3$
- 388 nm laser intensity: 0.5 W/cm²

Experiment result

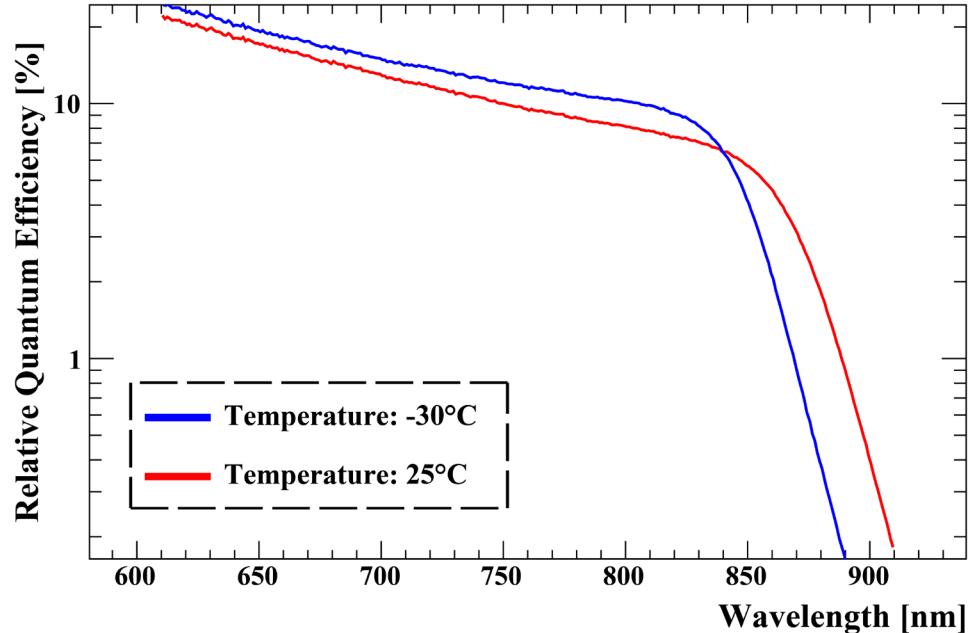
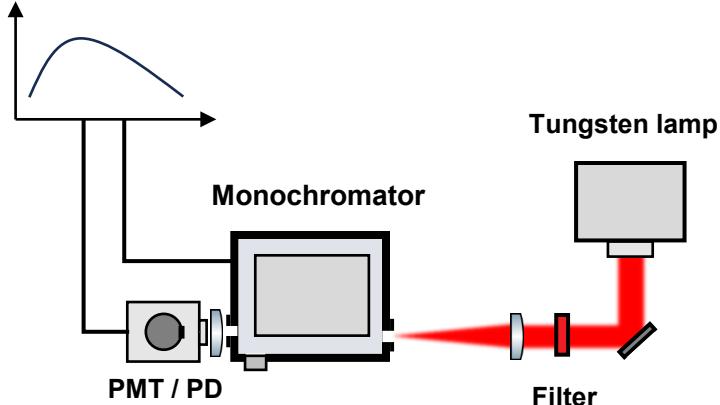
- Spectrums by scanning the monochromator



- Forbidden transitions are successfully observed
- Responsivity of the instruments is needed

PMT calibration

- Calibration experiment setup



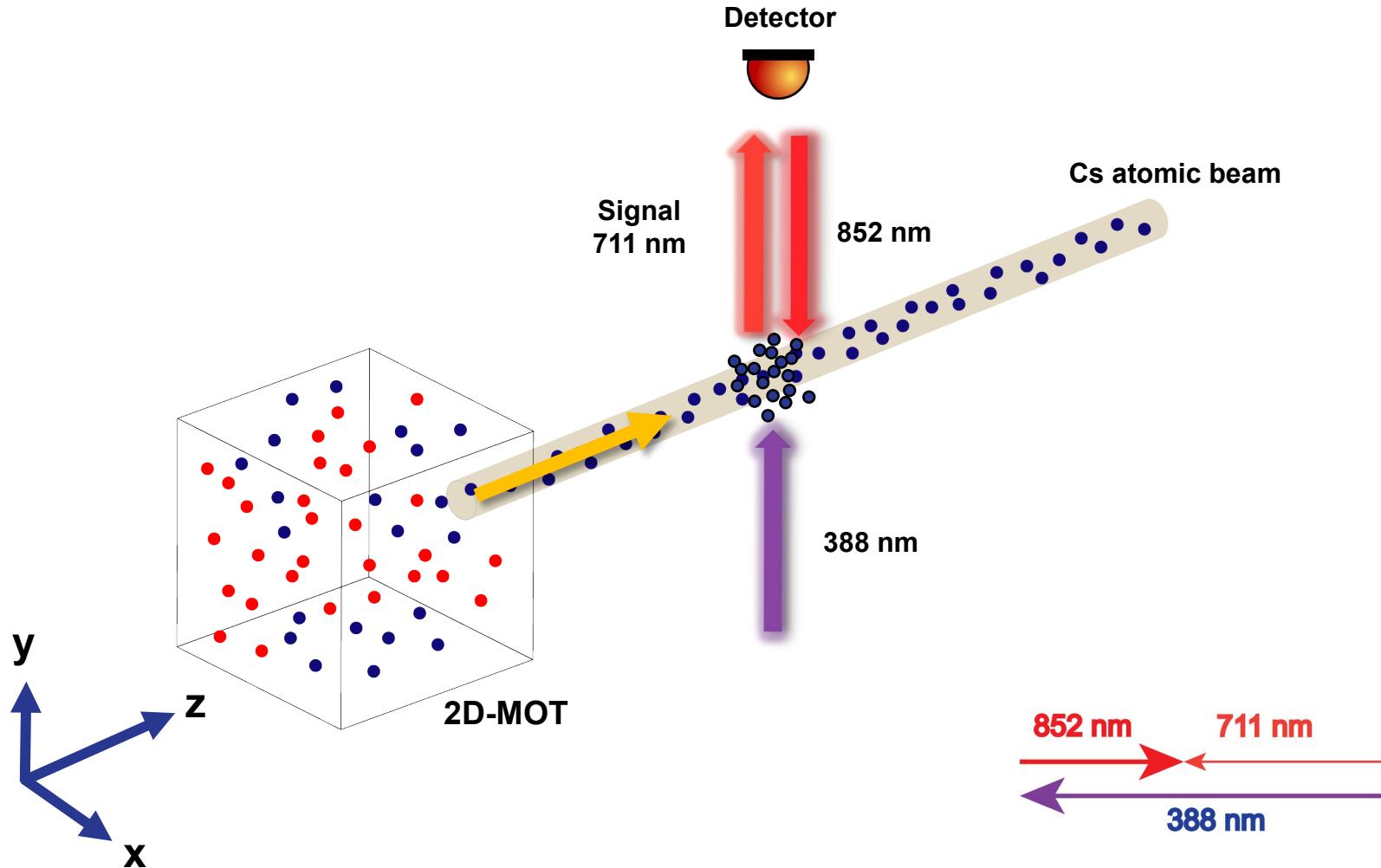
- Relative quantum efficiency of PMT is calibrated
- A coefficient of the forbidden transition is estimated (2.3 Hz for 711 nm)

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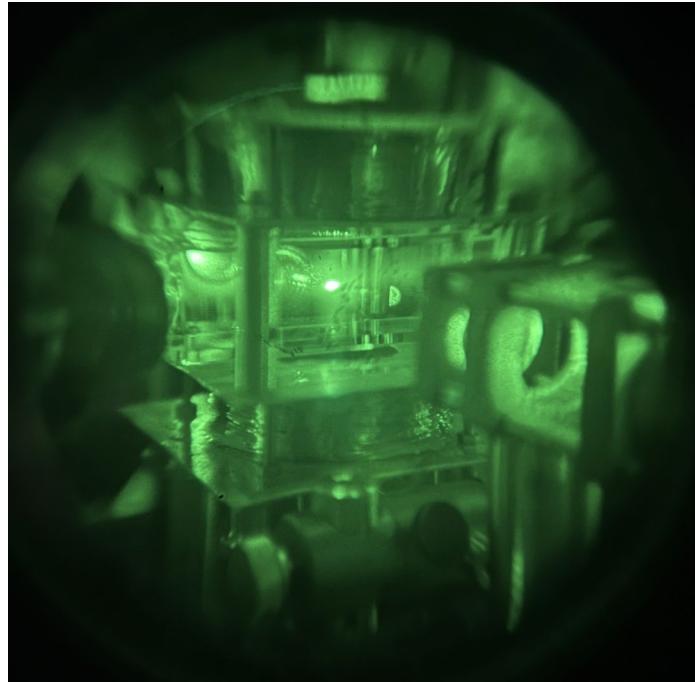
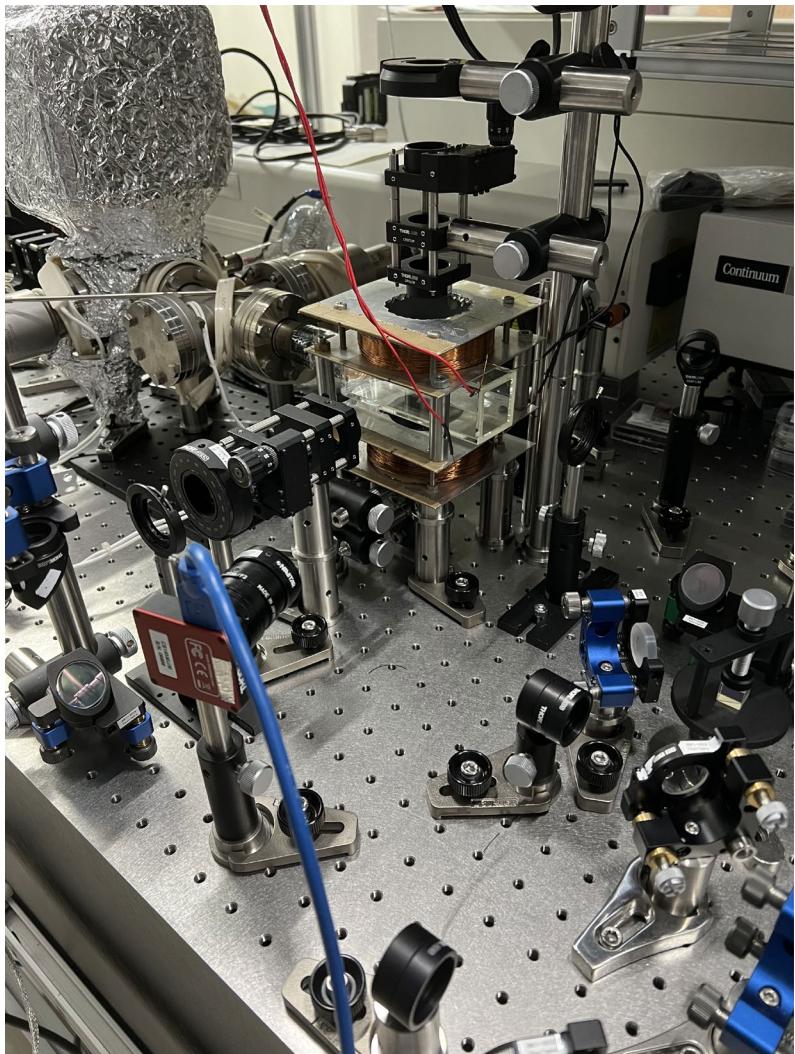
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Future Prospect

- Coherence measurement experiment using 2D Magneto-optical trap(MOT)



3D-MOT of Cesium



Summary

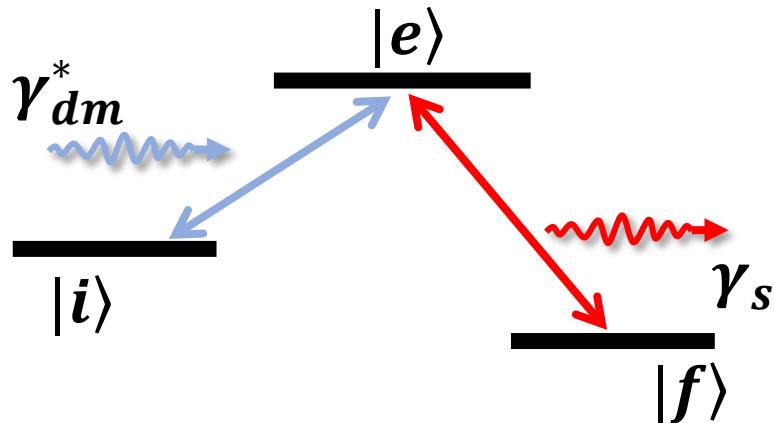
- ❑ New method of dark matter search.
 - ❖ Amplification by macroscopic coherence is KEY
- ❑ Coherence measurement experiment with Cs E2 transition
 - ✓ Determine E2 transition rate
 - ❖ Coherence measurement experiment
- ❑ Dark matter search in meV region

Thank you for your attention!

Mass scan

- Interaction Hamiltonian

$$\hat{H}_I = A \vec{L} \cdot \vec{S} + \mu_B (2\vec{S}_z + \vec{L}_z) \cdot \vec{B}_{ext}$$

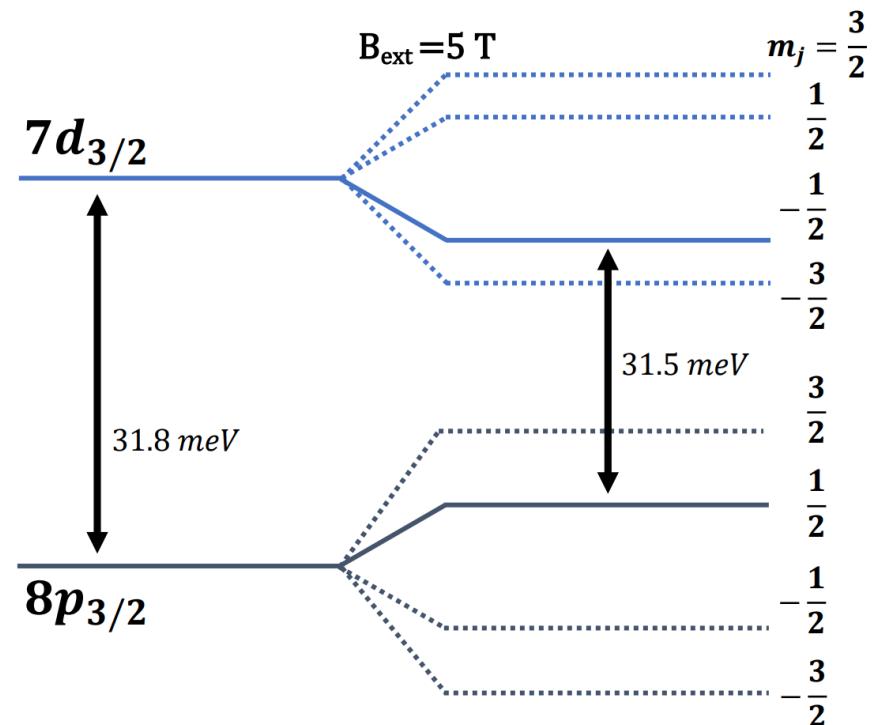


- Energy shift caused by Zeeman effect

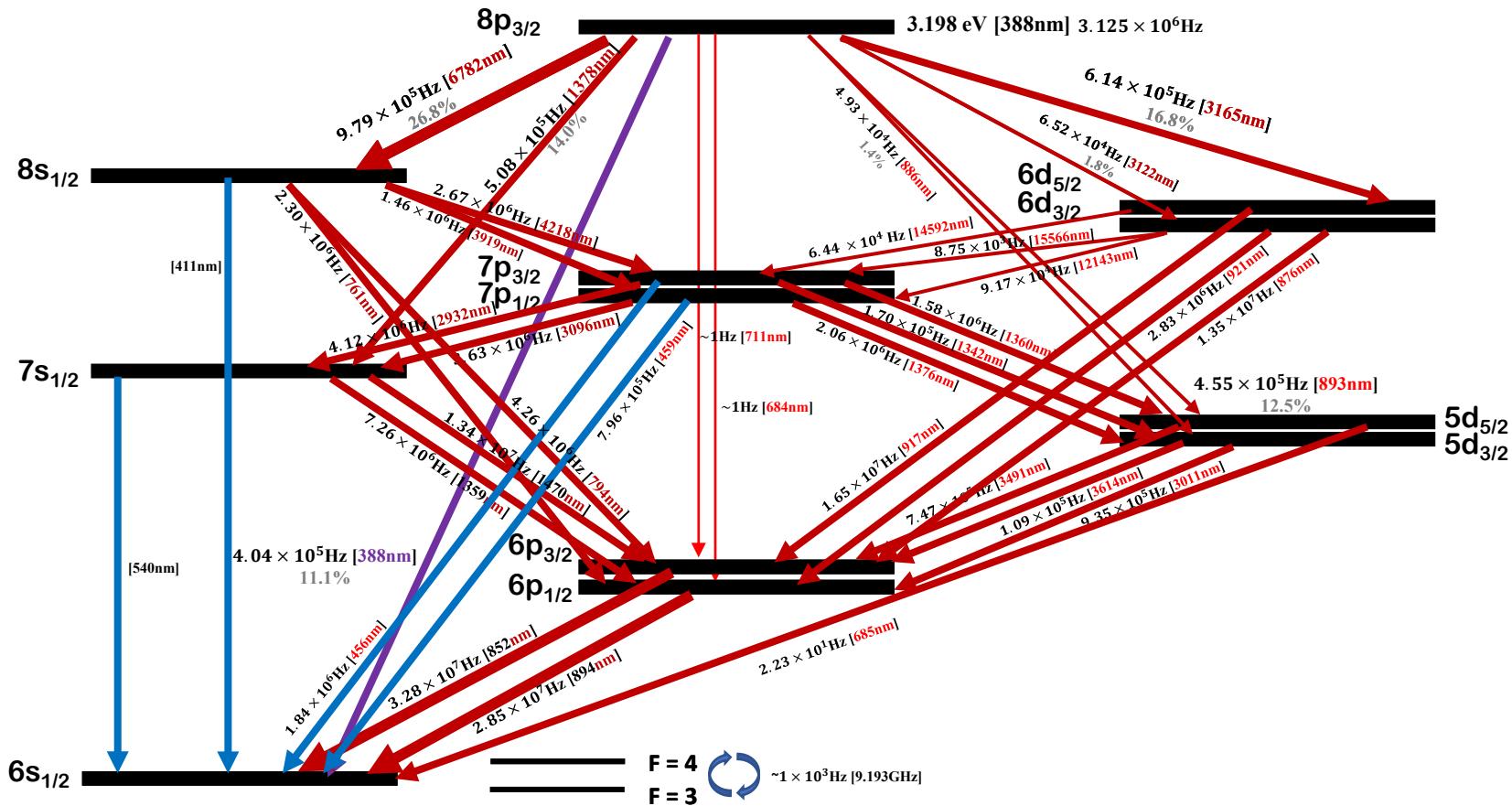
$$\Delta E_{m_j} = g_L \mu_B B_{ext} m_j$$

- $8p_{\{3/2, -1/2\}} \rightarrow 7d_{\{3/2, -1/2\}}$ under $5 \text{ T } B_{ext}$

$$\Delta m = -0.3 \text{ meV}$$



Cs Energy level



3D-MOT of Cesium

