

# Effect of Gravitational Cooling on Ultralight Dark Matter System Compared to Dynamical Friction

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(Collaborated with Jae-Weon Lee<sup>3</sup>, Dongsu Bak<sup>1,2</sup>)

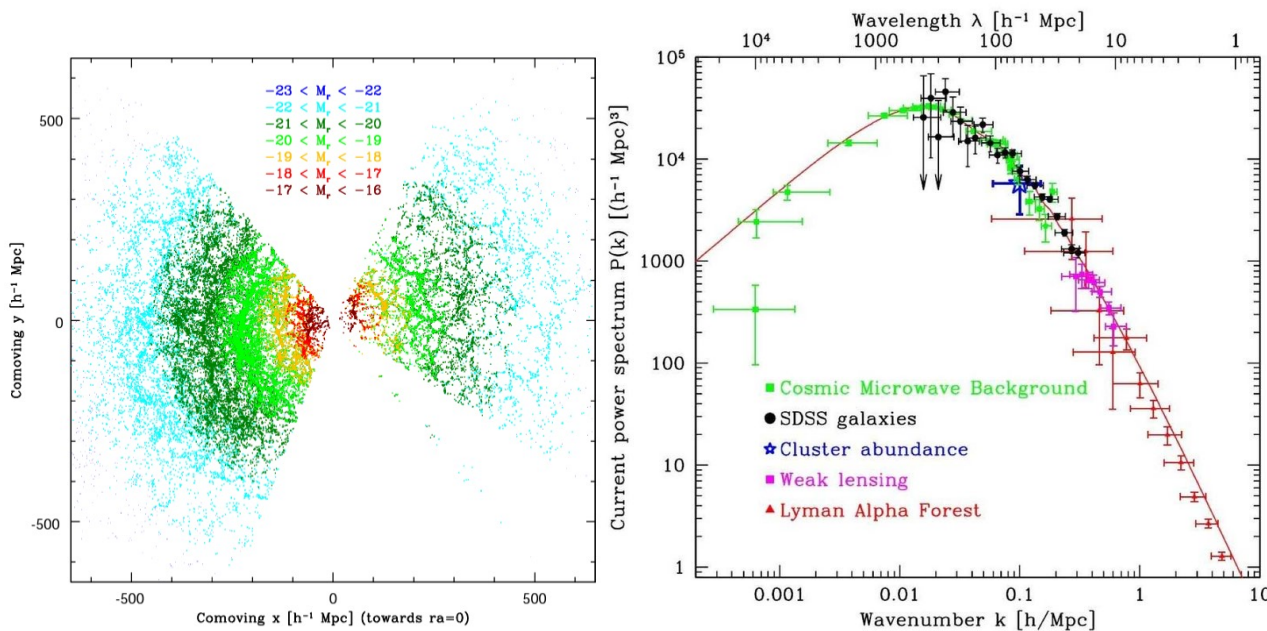
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# Does $\Lambda$ CDM explain our Universe well?

For Large Scale:  $\sim 0$ (Mpc)

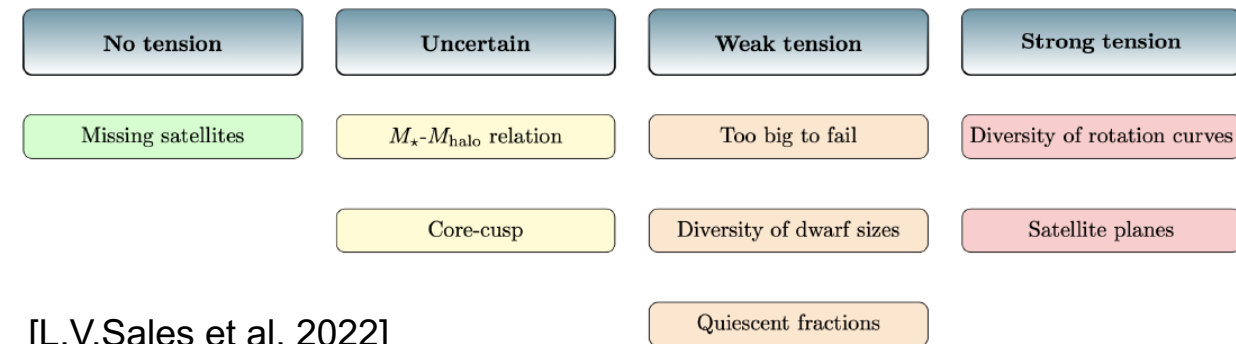


[M.Tegmark et al. 2003]

- $\Lambda$ CDM could explain the observed large-scale structure of our Universe well.

For Small Scale:  $\sim 0$ (kpc)

## $\Lambda$ CDM Tensions with Dwarf Galaxies



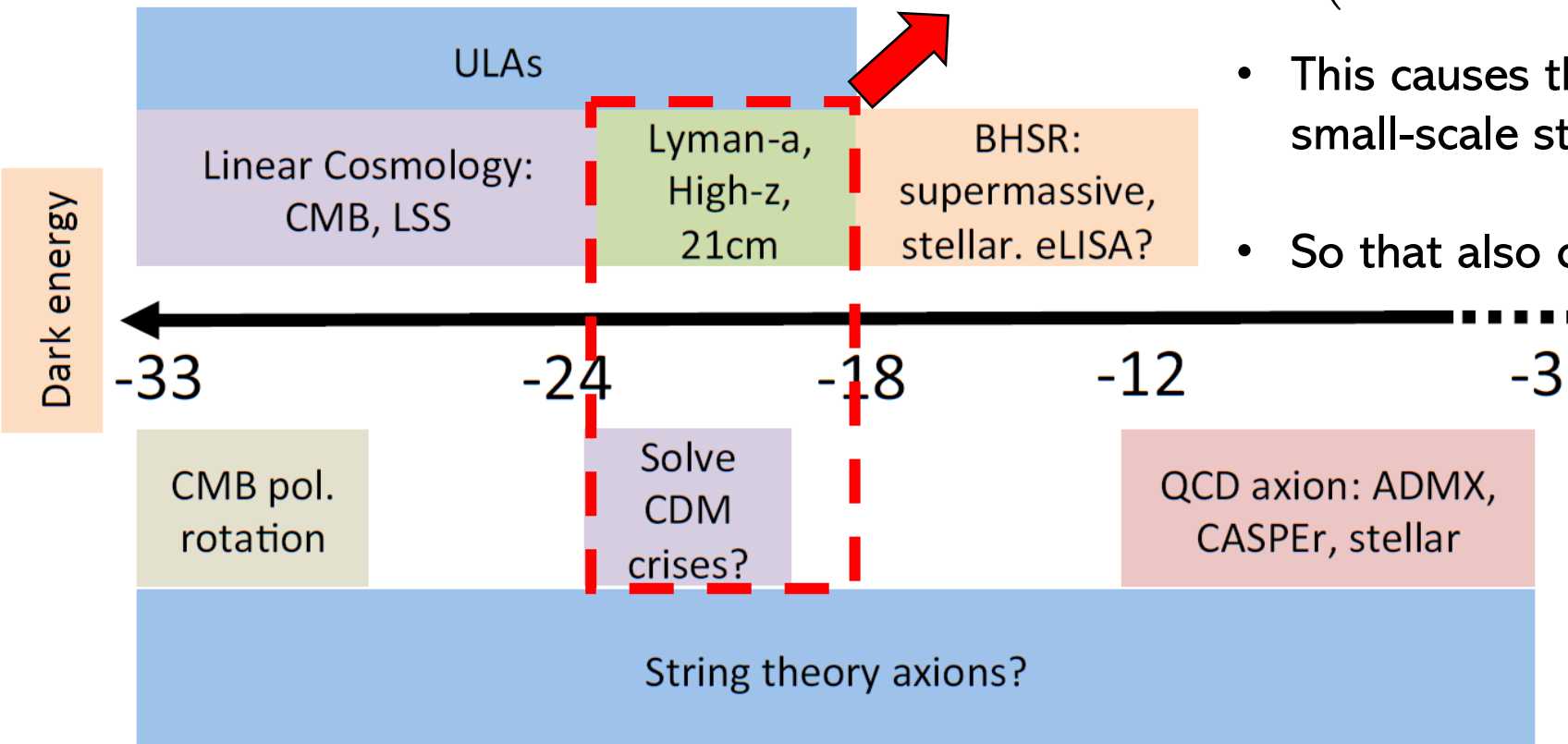
[L.V.Sales et al. 2022]

- Some problems can be solved by adding baryonic matter in  $\Lambda$ CDM simulation, but still not enough.
- **Alternative Dark Matter model is required to solve!**

# UltraLight Axion Dark Matter(ULDM)

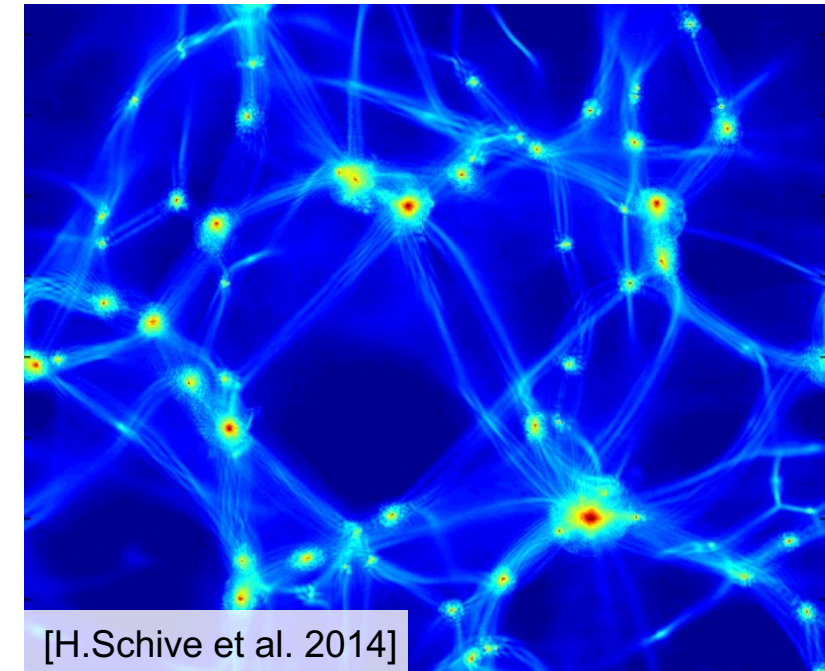
$$\lambda_{dB} = \frac{h}{m\sigma} \sim 1.204 \text{kpc} \left( \frac{10^{-22} \text{eV}/c^2}{m} \right) \left( \frac{100 \text{km/s}}{\sigma} \right)$$

- This causes the wave nature, suppressing the small-scale structure formation in our universe.
- So that also called “Fuzzy Dark Matter(FDM)”



$\log_{10}(m_a/\text{eV})$

[D.J.E.Marsh 2016]



[H.Schive et al. 2014]

$$\mathcal{L} = \frac{1}{2} g^{\mu\nu} (\partial_\mu \phi) (\partial_\nu \phi) - V(\phi), \quad V(\phi) = m^2 f_a^2 \left( 1 - \cos \frac{\phi}{f_a} \right) = \frac{1}{2} m^2 \phi^2 - \frac{\lambda}{4!} \phi^4 + \dots$$

$$\lambda = \left( \frac{m}{f_a} \right)^2$$

- The relic density of ALP today is:

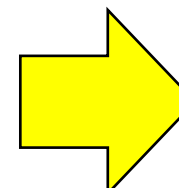
$$\Omega_{\text{axion}} \sim 0.1 \left( \frac{f_a}{10^{17} \text{ GeV}} \right)^2 \left( \frac{m}{10^{-22} \text{ eV}} \right)^{1/2} \quad [\text{L.Hui et al. 2016}]$$

All interaction terms of  $\sim f_a^{-1}$  are negligible!

This simplifies ALP system to  $V(\phi) \simeq \frac{1}{2} m^2 \phi^2$

Weak-Gravity Limit  $ds^2 = \left( 1 + \frac{2\Phi}{c^2} \right) c^2 dt^2 - \left( 1 - \frac{2\Phi}{c^2} \right) \eta_{ij} dx^i dx^j$

Nonrelativistic Limit  $\phi = \sqrt{\frac{\hbar^3 c}{2m}} \left( \psi e^{-i\frac{mc^2 t}{\hbar}} + \psi^* e^{i\frac{mc^2 t}{\hbar}} \right)$



[Schrodinger-Poisson Eq.]

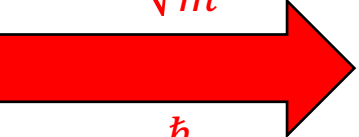
$$-\frac{\hbar^2}{2m} \nabla^2 \psi + m\Phi\psi = i\hbar \frac{\partial \psi}{\partial t}$$

$$\nabla^2 \Phi = 4\pi G m |\psi|^2$$


# Madelung Formalism of SP eq.

$$-\frac{\hbar^2}{2m}\nabla^2\psi + m\Phi\psi = i\hbar\frac{\partial\psi}{\partial t}$$

$\psi = \sqrt{\frac{\rho}{m}}e^{i\theta}$   
 $\vec{v} = \frac{\hbar}{m}\nabla\theta$



$$\left\{ \begin{array}{l} 1. \text{ Continuity eq. } \frac{\partial\rho}{\partial t} + \nabla \cdot (\rho\vec{v}) = 0 \\ 2. \text{ Euler eq. } \left(\frac{\partial}{\partial t} + \vec{v} \cdot \nabla\right)\vec{v} = -\nabla\Phi - \nabla Q \end{array} \right.$$



$$3. \text{ Poisson eq. } \nabla^2\Phi = 4\pi G\rho$$

- $Q$  roles the quantum pressure, which resists to the gravitational collapse of ULDM system.

$$Q = -\frac{\hbar^2}{2m^2}\left(\frac{\nabla^2\sqrt{\rho}}{\sqrt{\rho}}\right)$$

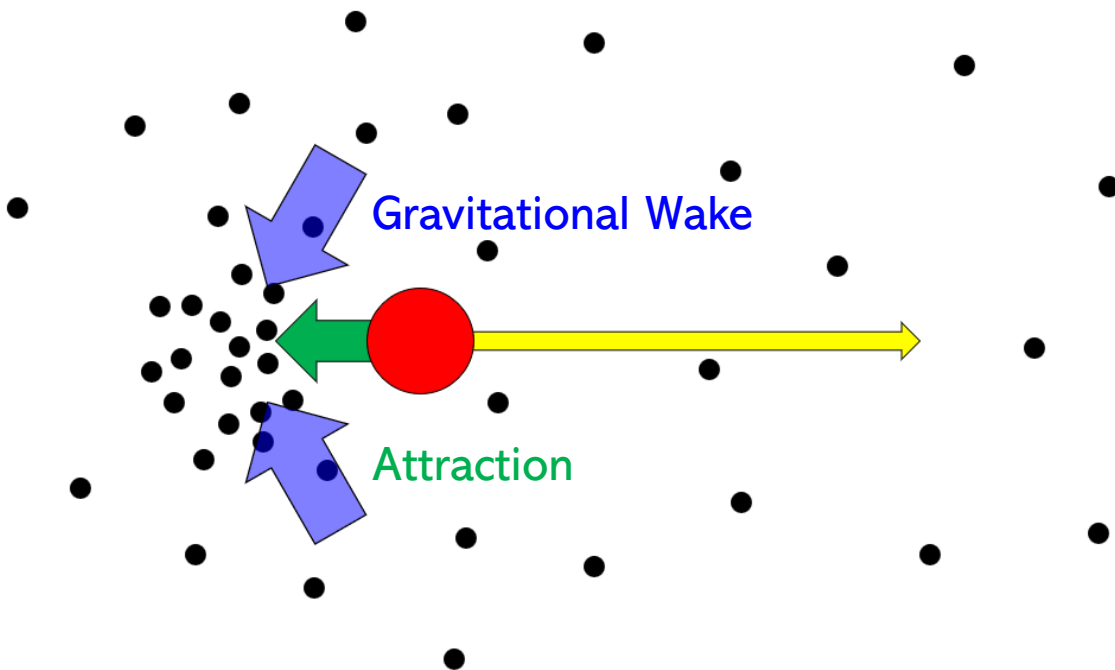
# Dynamical Friction of CDM/ULDM

- An object of mass  $M$  traveling inside a particle field of density  $\bar{\rho}$  with velocity  $v$  experiences the dynamical friction as below:

$$F_{\text{DF}} = 4\pi\bar{\rho} \left( \frac{GM}{v} \right)^2 \boxed{C_{\text{DF}}}$$

DF Coefficient!

## 1. For Classical Particles (CDM)

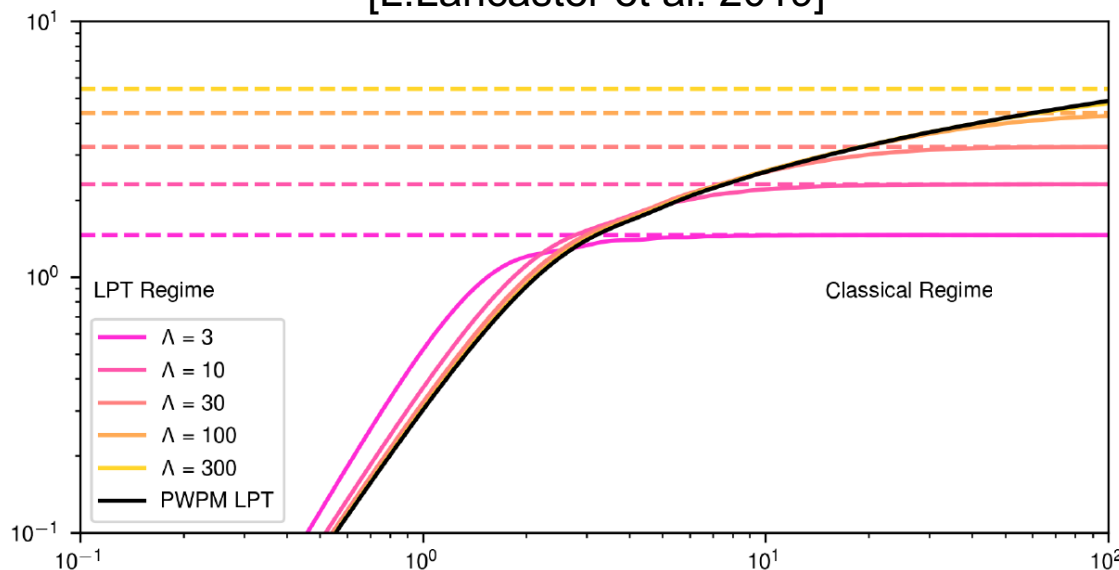
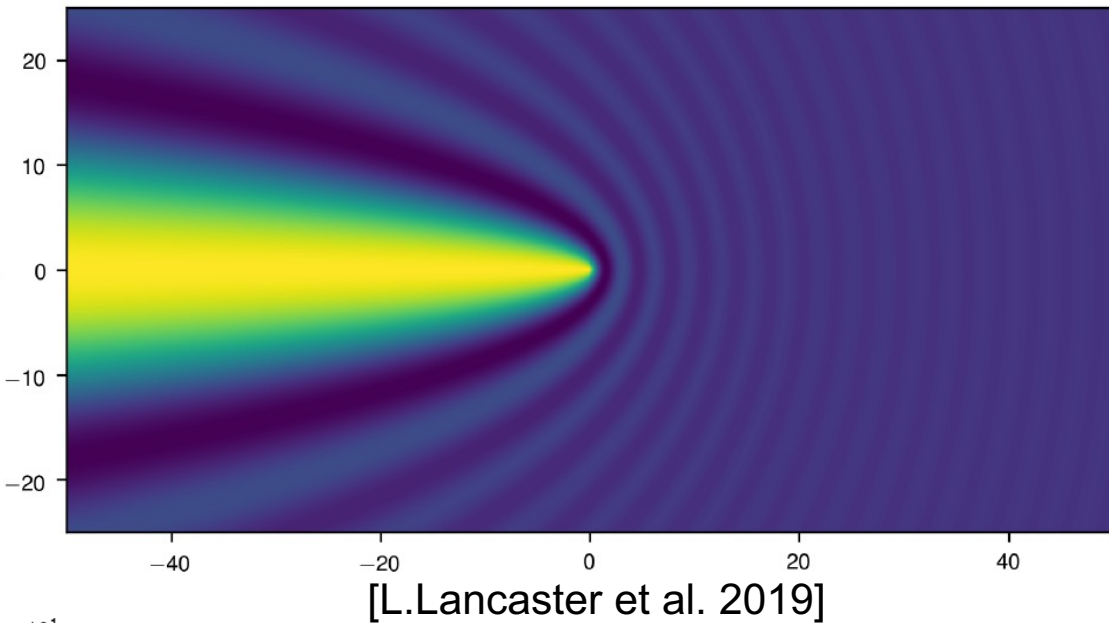


- The object loses its kinetic energy by gravitational interactions with surrounding field particles. [Chandrasekhar 1943]
- For sufficiently large  $v$ ...

$$\frac{d\vec{v}_M}{dt} = -4\pi G^2 M m_a n \log \Lambda \frac{\vec{v}}{v^3} \quad [\text{Binney \& Tremaine}]$$

$$C_{\text{DF}} = \log \Lambda = \log \frac{b_{\text{max}}}{\max(r_h, GM/v^2)}$$

## 2. For Ultralight Dark Matter



- Linear Perturbation of Madelung eq. explains the dynamical friction of ULDM well.
- $C_{DF}$  also can be derived from exact solution of SP eq.

$$C_{DF}(\tilde{b}, \beta) \simeq \begin{cases} \frac{1}{3} \tilde{b}^2 & (\tilde{b} \ll 1) \\ \log(2\tilde{b}) - 1 + \text{Re}\Psi(1 + i\beta) & (\tilde{b} \gg 1) \end{cases}$$

$\tilde{b} \equiv \frac{m v b}{\hbar}$        $\beta = \frac{G M m}{\hbar v}$       [L.Hui et al. 2016]

- For  $\tilde{b} \gg 1$ , the dynamical friction of ULDM becomes similar to that of CDM(classical limit).
- **Dynamical friction of ULDM is generally less than of CDM.**



# Gravitational Cooling Effect of ULDM

# What is Gravitational Cooling Effect of ULDM?

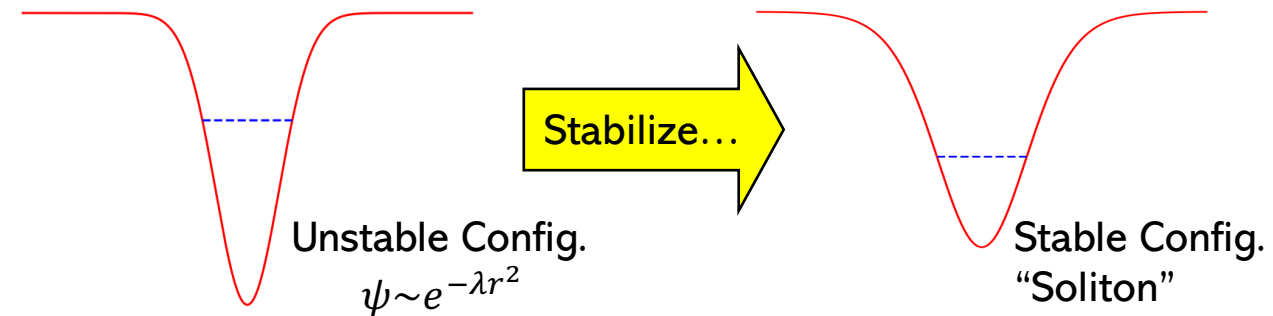
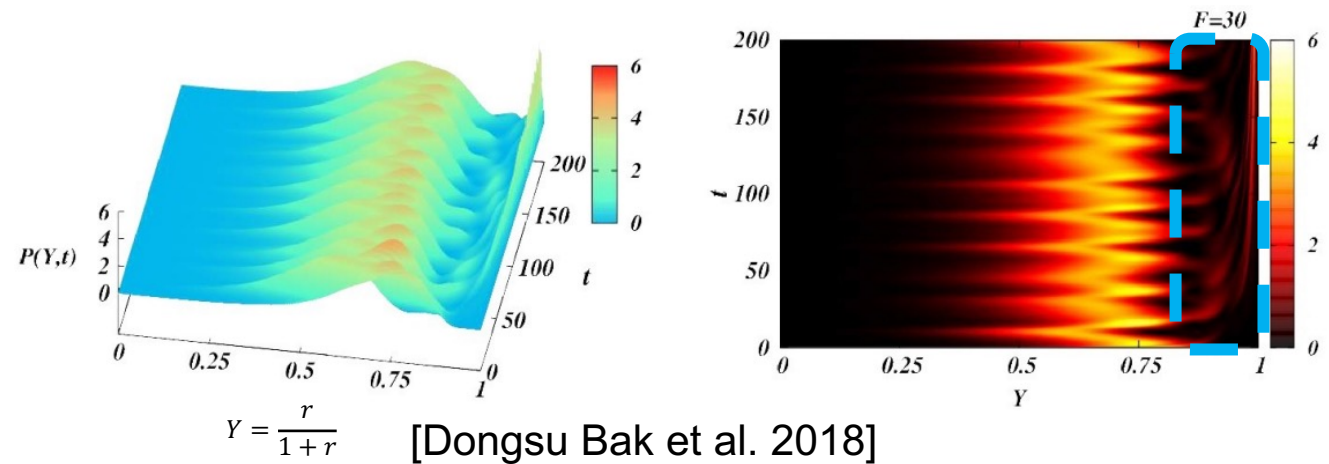
- A relaxation process of the unstable configuration, through the emission of DM particles. The wave nature allows high- $k$  modes with high energy  $\epsilon_k = \frac{\hbar^2 k^2}{2m}$  to escape from the potential well.

- Its timescale  $\tau_{gc}$  has a relation with total mass of system as:

$$\tau_{gc} = \left(\frac{\hbar}{m}\right)^3 \frac{1}{(GM_{\text{tot}})^2}$$

$$= 0.372 \text{ Myr} \left(\frac{10^{-22} \text{ eV}}{m}\right)^3 \left(\frac{10^9 M_{\odot}}{M_{\text{tot}}}\right)^2$$

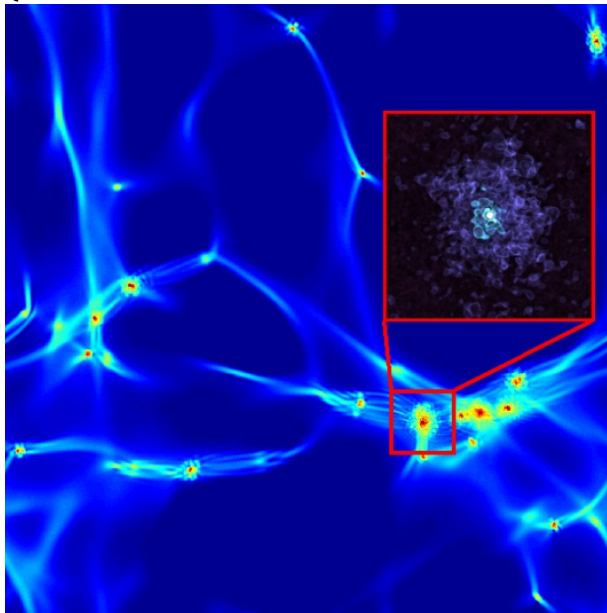
- This effect had only been derived from numerical calculation yet.



# Solitonic Core and ULDM Halo

- The most stable ground-state solution of SP eq.

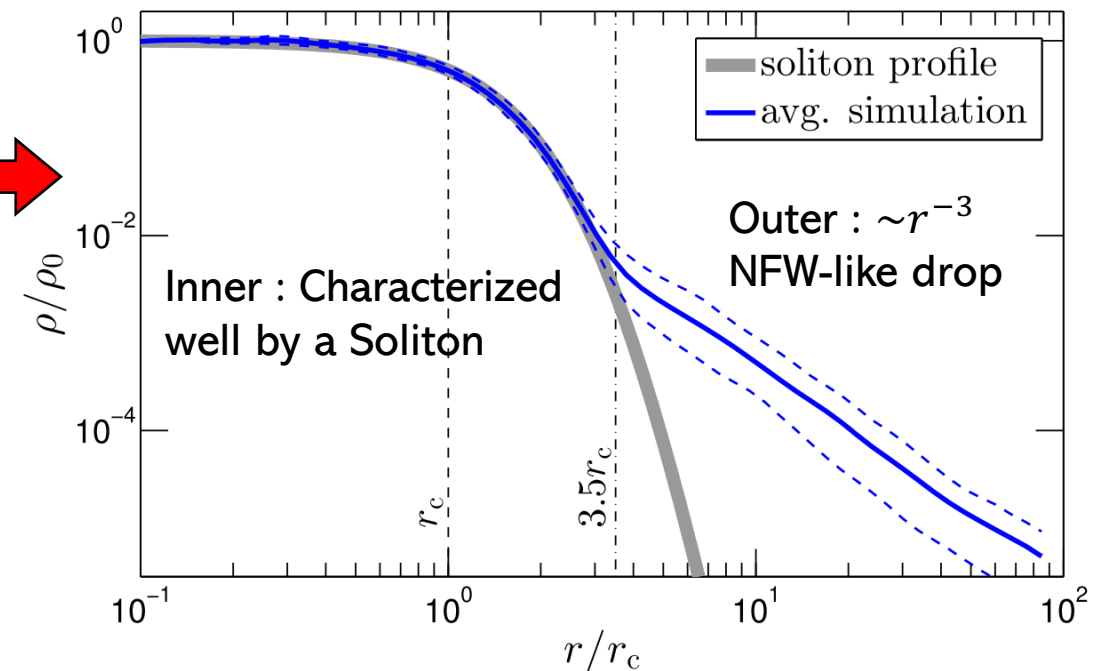
$$\left\{ \begin{array}{l} r_{1/2} \cong 0.335 \text{kpc} \frac{10^9 M_{\odot}}{M} \left( \frac{10^{-22} \text{eV}}{m} \right)^2 \\ \rho_c \cong 7.05 M_{\odot} / \text{pc}^3 \left( \frac{m}{10^{-22} \text{eV}} \right)^6 \left( \frac{M}{10^9 M_{\odot}} \right)^4 \end{array} \right.$$



[P.Mocz et al. 2017]

$$\rho_{\text{soliton}}(r) \cong \frac{\rho_c}{\{1 + 0.091(r/r_c)^2\}^8} \quad [\text{L.Hui et al. 2016}]$$

$r_h \cong 1.452 r_c$



# Some Results of Numerical Simulations

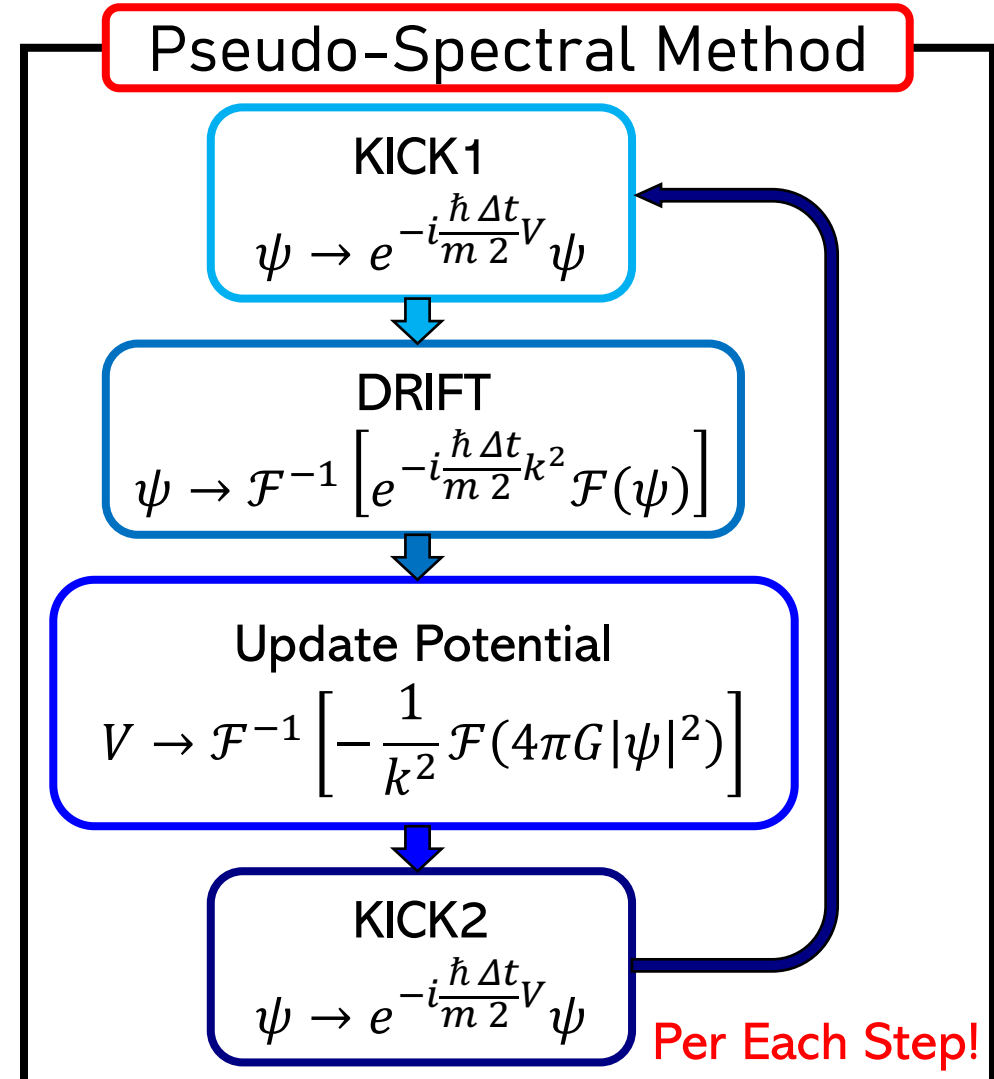
# Schrodinger-Poisson eq. Solver

- PyUltraLight : Code for solving SP equation using Pseudospectral Method with periodic boundary condition [F.Edwards et.al. 2018]
- PyUltraLight2 : SP equation coupled with N-body particles solved by 4<sup>th</sup> order Runge-Kutta method [W.Wang et.al. 2021]

$$\left\{ \begin{array}{l} -\frac{\hbar^2}{2m} \nabla^2 \psi + m(\Phi_U + \Phi_N) \psi = i\hbar \frac{\partial \psi}{\partial t} \\ \nabla^2 \Phi_U = 4\pi G |\psi|^2 \end{array} \right. \oplus \left\{ \begin{array}{l} \Phi_N(\vec{x}) = -\sum_j \frac{GM_j}{|\vec{x} - \vec{x}_j|} \\ \frac{d^2 \vec{x}_j}{dt^2} = -\nabla_j \sum_{k \neq j} \frac{GM_k}{|\vec{x}_j - \vec{x}_k|} - \nabla \Phi_U(\vec{x}_j) \end{array} \right.$$

- Timestep is chosen to prevent the artifact from being canceled out by constructive interference of  $\psi$  when  $\arg(\psi) \geq \pi$ .

$$\Delta t = \frac{m \Delta x^2}{\hbar \pi}$$



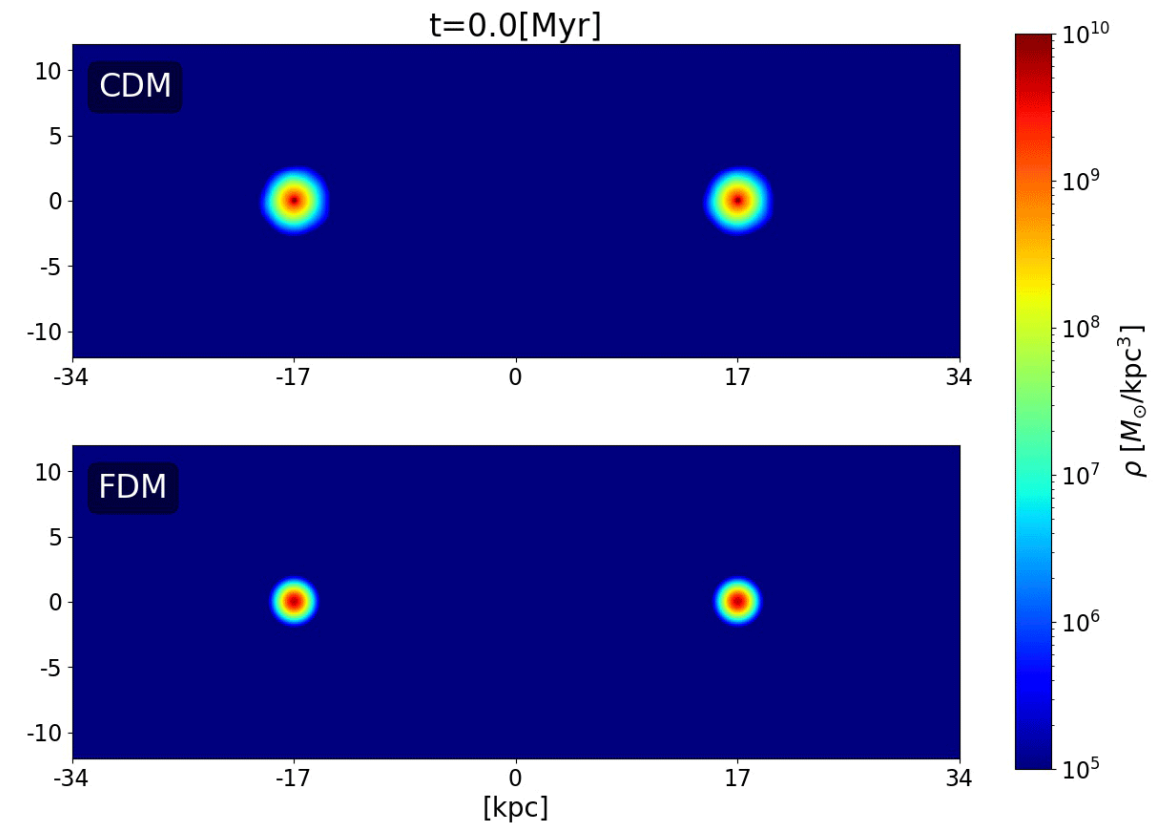
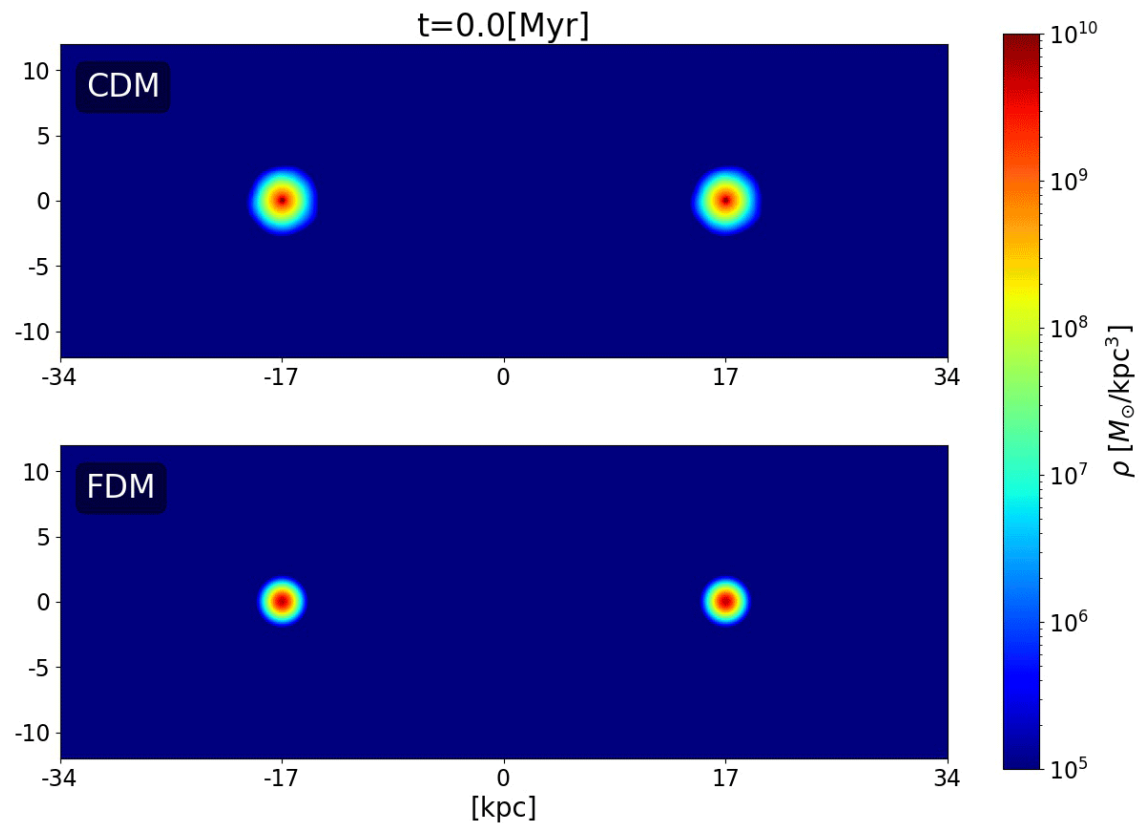
# 1. Head-on Collision of same-mass Halos

[H.Koo et al. 23XX.XXXXX]

- Each CDM/ULDM(FDM) halo has a mass of  $M = 2\pi \times 10^8 M_\odot$  & ULDM particle mass  $m = 10^{-22} \text{eV}/c^2$

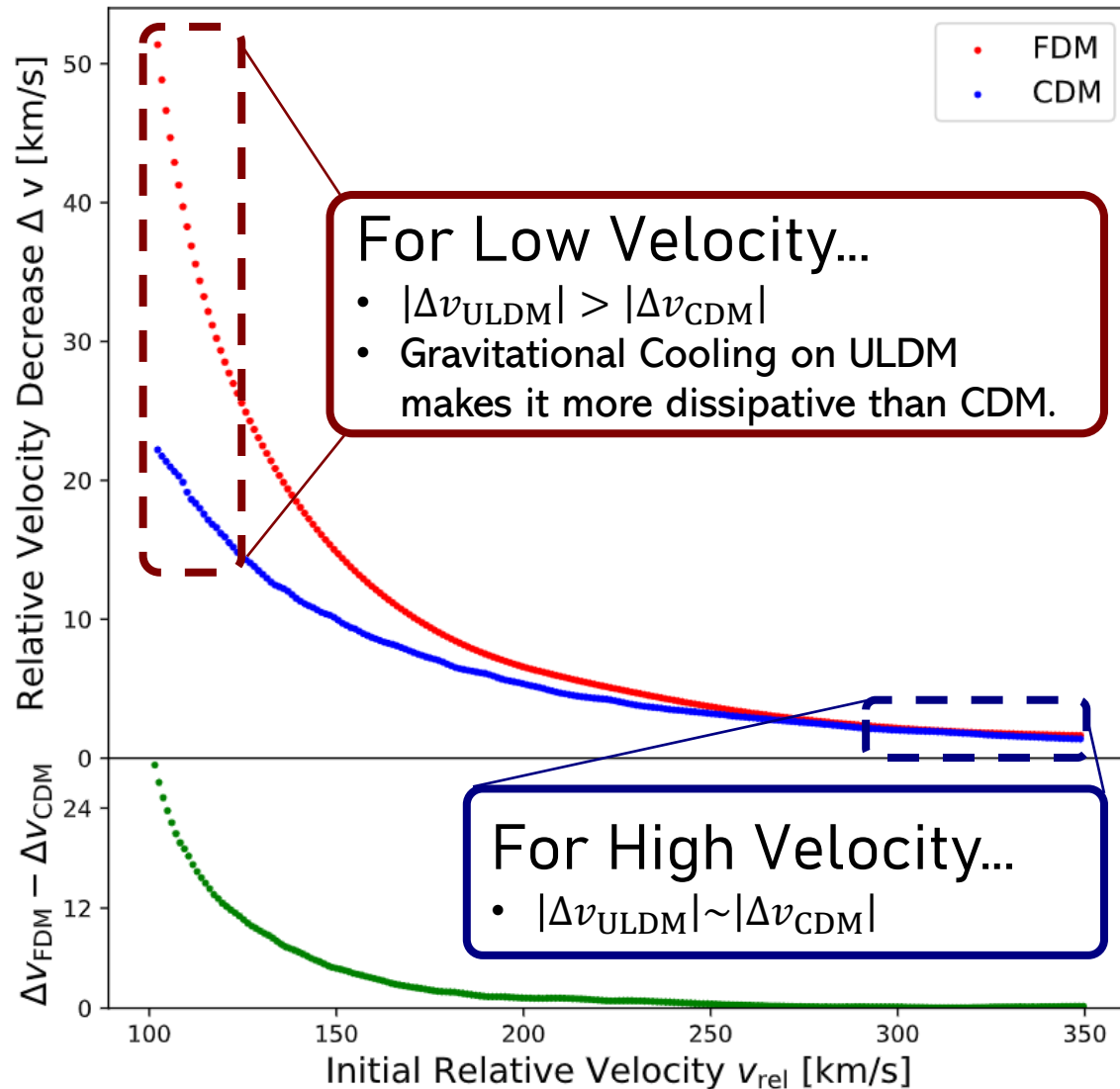
$$v_0^{\text{rel}} = 112.78 \text{ km/s}$$

$$v_0^{\text{rel}} = 225.56 \text{ km/s}$$



# 1. Head-on Collision of same-mass Halos

[H.Koo et al. 23XX.XXXXXX]



$$q \equiv \frac{\hbar}{mvr_{1/2}}$$

$v = 100 \sim 350 \text{ km/s} \rightarrow q = 0.103 \sim 0.360$

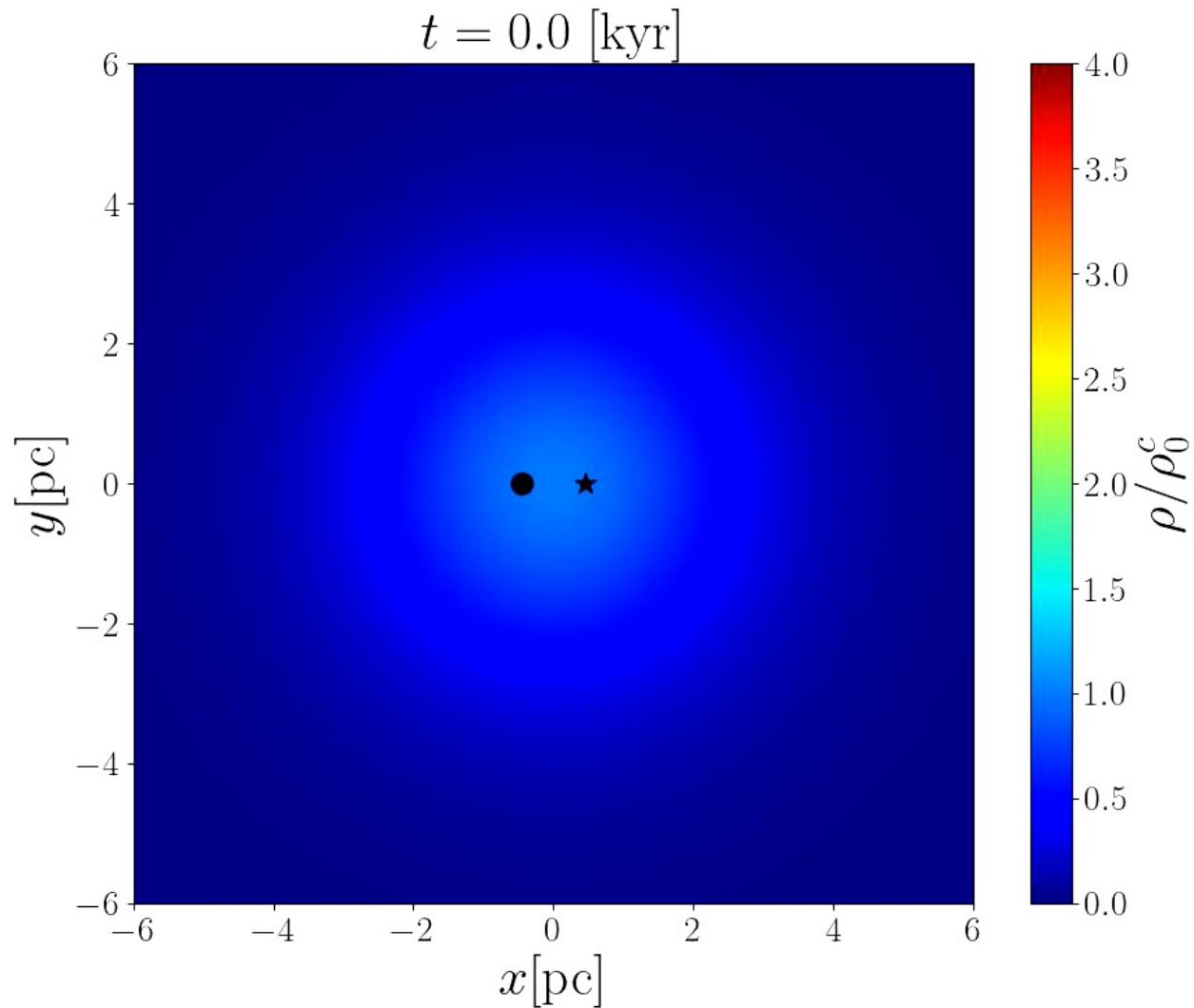
$$\left| \frac{\Delta v}{v} \right|_{\text{CDM}} = A_c q^2 (-\log q + B_c) \begin{cases} A_c = 15.93 \pm 0.126 \\ B_c = -0.169 \pm 0.008 \end{cases}$$

$$\left| \frac{\Delta v}{v} \right|_{\text{ULDM}} = A_u q^3 (1 + B_u q^2) \begin{cases} A_u = 3.692 \pm 0.041 \\ B_u = 15.63 \pm 0.294 \end{cases}$$

Fitting Function Derived from [D.Bak et al. 2010.14738]

## 2. Supermassive BH Binary in ULDM Halo

[H.Koo et al. 2311.03412]



$$m = 10^{-21} \text{eV}/c^2$$

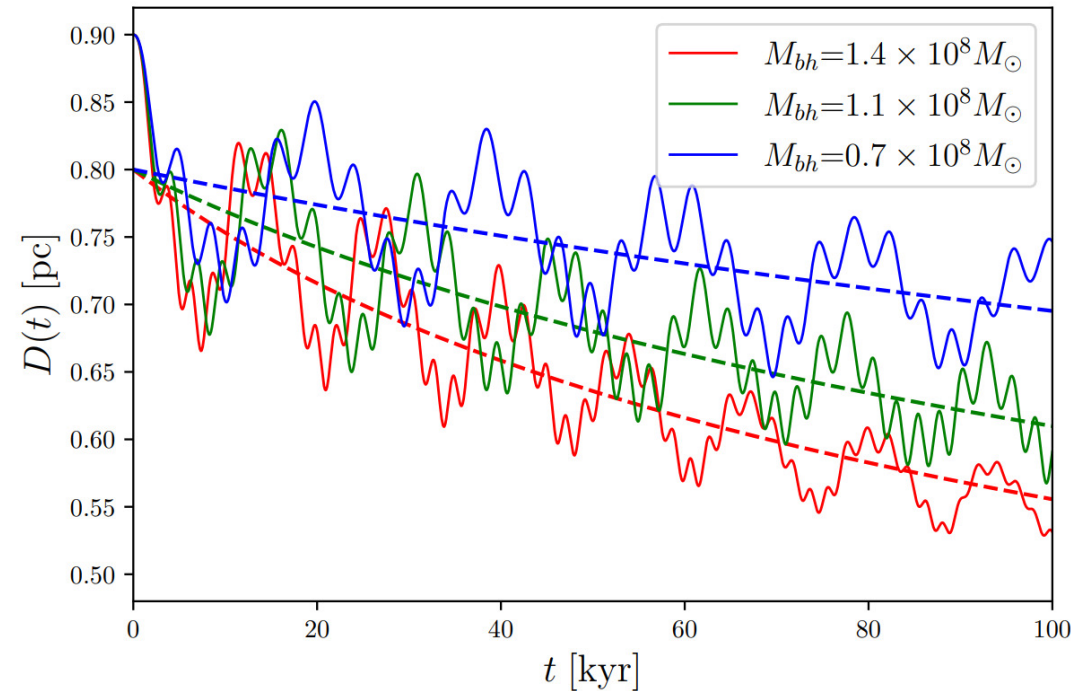
$$M_s = 10^9 M_\odot$$

$$M_{\text{bh}} = 10^8 M_\odot$$



## 2. Supermassive BH Binary in ULDM Halo

[H.Koo et al. 2311.03412]



Gravitational Cooling on total system could be

1. Global Mode : Breathing soliton with period  $\tau$

$$\tau = \tau_0 \left( \frac{10^{-21} \text{eV}}{m} \right)^3 \left( \frac{10^9 M_\odot}{M_s} \right)^2 \left( 1 + 2\gamma_g \frac{M_{\text{bh}}}{M_s} \right)^{-2}$$

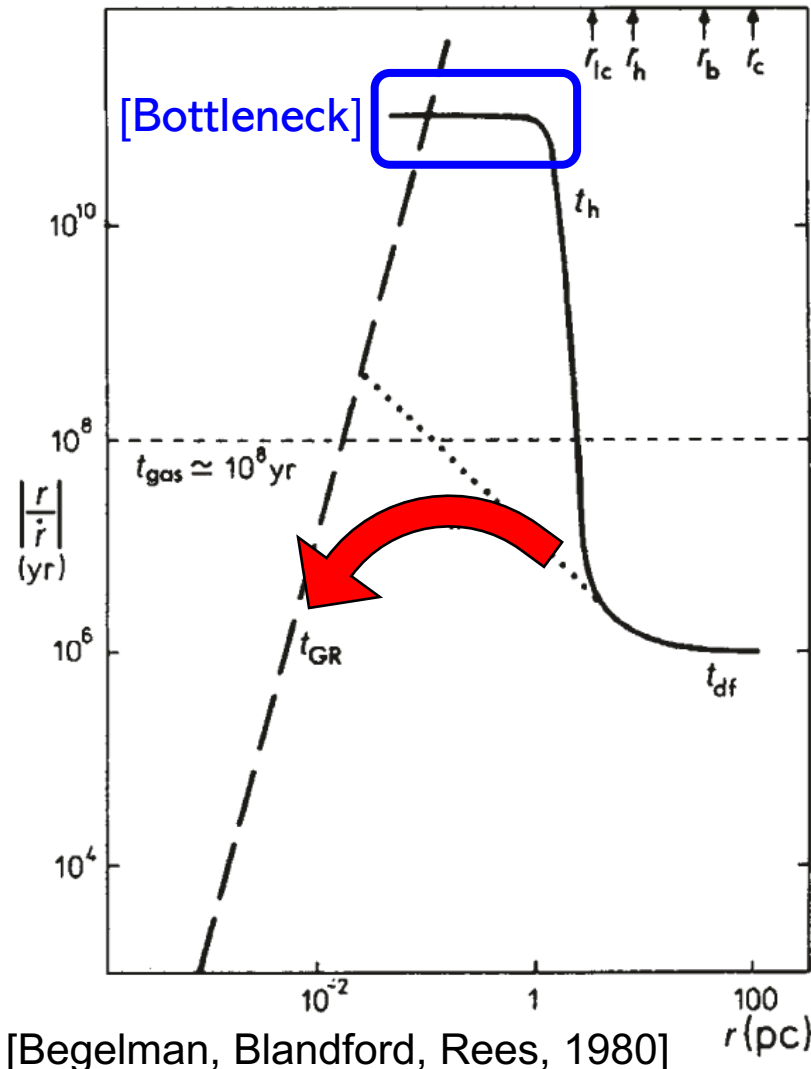
$$\begin{cases} \gamma_g = 1.974 \pm 0.013 \cong 2 \\ \tau_0 = (32.11 \pm 0.115) \text{kyr} \end{cases}$$

2. Local Mode : Orbital Decay of SMBHB

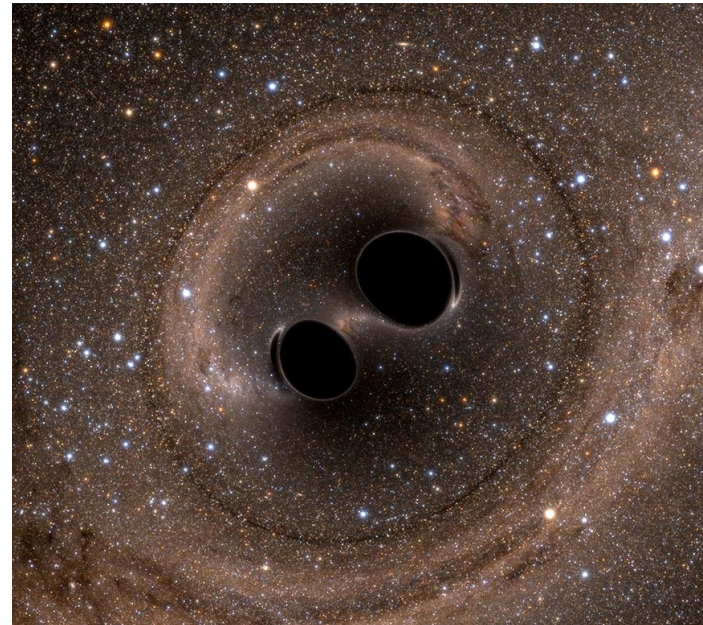
$$\bar{D}(t) = \bar{D}_0 (1 + 3Q\bar{D}_0^3 t)^{-\frac{1}{3}}$$

$$Q = \kappa \left( \frac{m}{10^{-21} \text{eV}/c^2} \right)^9 \frac{(M_s + 2\gamma_l M_{\text{bh}})^6}{M_{\text{bh}}} \begin{cases} \kappa = \frac{29.58 \pm 1.39}{[\text{kyr} \cdot \text{pc}^3 \cdot (10^{10} M_\odot)^5]} \\ \gamma_l = 1.476 \pm 0.067 \cong 1.5 \end{cases}$$

# Can ULDM solve the Final Parsec Problem?



- Final Parsec Problem : Numerical simulations with  $\Lambda$ CDM show that SMBHBs typically stall out at  $10^{-2}\text{pc} < r < 1\text{pc}$  and take longer than the age of the universe to reach the era of gravitational radiation.



- For  $10^5 M_{\odot} \leq M_{\text{bh}} \leq 10^7 M_{\odot}$  inside ULDM soliton, the Gravitational Cooling Effect gives a hint for SMBHBs to survive from the timescale-bottleneck!

# Summary

- Ultralight Axion Dark Matter of  $m \sim 10^{-22} \text{eV}/c^2$  and  $\lambda_{\text{dB}} \sim O(\text{kpc})$  could solve some problems that  $\Lambda\text{CDM}$  encounters.
- The Gravitational Cooling Effect of ULDM system is a sort of relaxation through the emission of DM particles from potential well. This is significant for ULDM halo dynamics, rather than the dynamical friction.
- The Gravitational Cooling also could give a hint to solve Final Parsec Problem.
- Adding self interaction  $\mathcal{L}_I \sim \lambda \phi^4$  is expected to give more precise constraint!

# Appendix 1 : $C_{DF}$ for ULDM

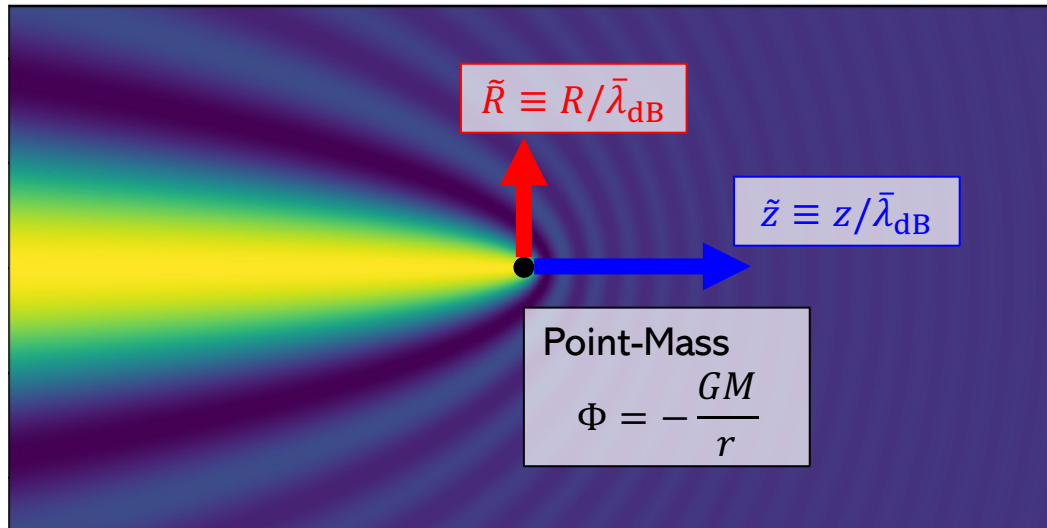
## 1. From Linear Perturbation

$$\begin{cases} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0 \\ \left( \frac{\partial}{\partial t} + \vec{v} \cdot \nabla \right) \vec{v} = -\nabla \Phi - \nabla Q \\ \nabla^2 \Phi = 4\pi G \rho \end{cases}$$



$$\begin{aligned} \rho &= \bar{\rho}(1 + \alpha) \\ \vec{v} &= v_0 \hat{z} + \delta \vec{v} \\ \psi_0(\vec{x}) &= \sqrt{\bar{\rho}} e^{i \frac{mv_0 z}{\hbar}} \end{aligned}$$

$$\begin{aligned} \left( \frac{\partial}{\partial t} - v_0 \frac{\partial}{\partial z} \right)^2 \alpha - \nabla^2 \Phi + \frac{\hbar^2}{4m^2} \nabla^4 \alpha &= 0 \\ \alpha(\tilde{R}, \tilde{z}) &= \frac{4\beta}{\pi} \int_0^\infty d\tilde{k}_R \int_{-\infty}^\infty d\tilde{k}_z \frac{\tilde{k}_R J_0(\tilde{k}_R \tilde{R}) e^{i\tilde{k}_z \tilde{z}}}{\tilde{k}^4 - 4\tilde{k}_z^2} \end{aligned}$$



- Dynamical friction from overdensity  $F_{DF} = \bar{\rho} \int d^3 \vec{x} \alpha(\vec{x}) (\hat{x}_{\parallel} \cdot \nabla) \Phi(\vec{x})$  leads to  $C_{DF}$  as:

$$C_{DF}(\tilde{b}) = - \int_0^{\tilde{b}} d\tilde{R} \int_0^{\sqrt{\tilde{b}^2 - \tilde{R}^2}} d\tilde{z} \int_0^2 du \frac{\tilde{R} \tilde{z} J_0(\sqrt{2u - u^2} \tilde{R}) \sin \tilde{z} u}{(\tilde{R}^2 + \tilde{z}^2)^{\frac{3}{2}} u}$$

# Appendix 1 : $C_{\text{DF}}$ for ULDM

2. From direct solution of SP eq. with Point-Mass

$$\psi(\tilde{R}, \tilde{z}) = \sqrt{\bar{\rho}} e^{i\tilde{z} + \frac{\pi}{2}\beta} |\Gamma(1 - i\beta)| \times M \left[ i\beta, 1; i \left( \sqrt{\tilde{R}^2 + \tilde{z}^2} + \tilde{z} \right) \right]$$

Where  $M$  is the confluent hypergeometric function.

$$\alpha(\vec{x}) \equiv \frac{\rho(\vec{x})}{\bar{\rho}} - 1 = \frac{|\psi|^2}{\bar{\rho}} - 1$$

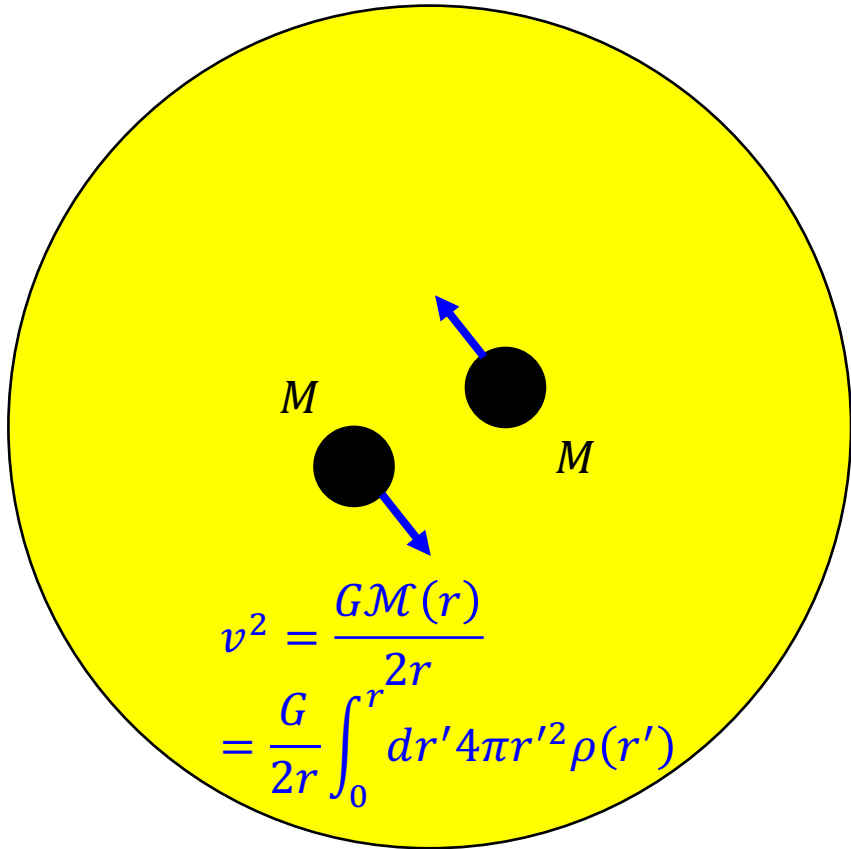
We samely put  $\alpha$  into  $F_{\text{DF}} = \bar{\rho} \int d^3\vec{x} \alpha(\vec{x}) (\hat{x}_{\parallel} \cdot \nabla) \Phi(\vec{x})$ , then could get:

$$C_{\text{DF}}(\tilde{b}, \beta) = \frac{1}{2\beta} e^{\pi\beta} |\Gamma(1 - i\beta)|^2 \int_0^{2\tilde{b}} d\tilde{q} |M(i\beta, 1; i\tilde{q})|^2 \left( \frac{\tilde{q}}{\tilde{b}} - 2 - \log \frac{\tilde{q}}{2\tilde{b}} \right)$$

$$\simeq \begin{cases} \frac{1}{3} \tilde{b}^2 (\tilde{b} \ll 1) \\ \log(2\tilde{b}) - 1 + \text{Re}\Psi(1 + i\beta) (\tilde{b} \gg 1) \end{cases}$$

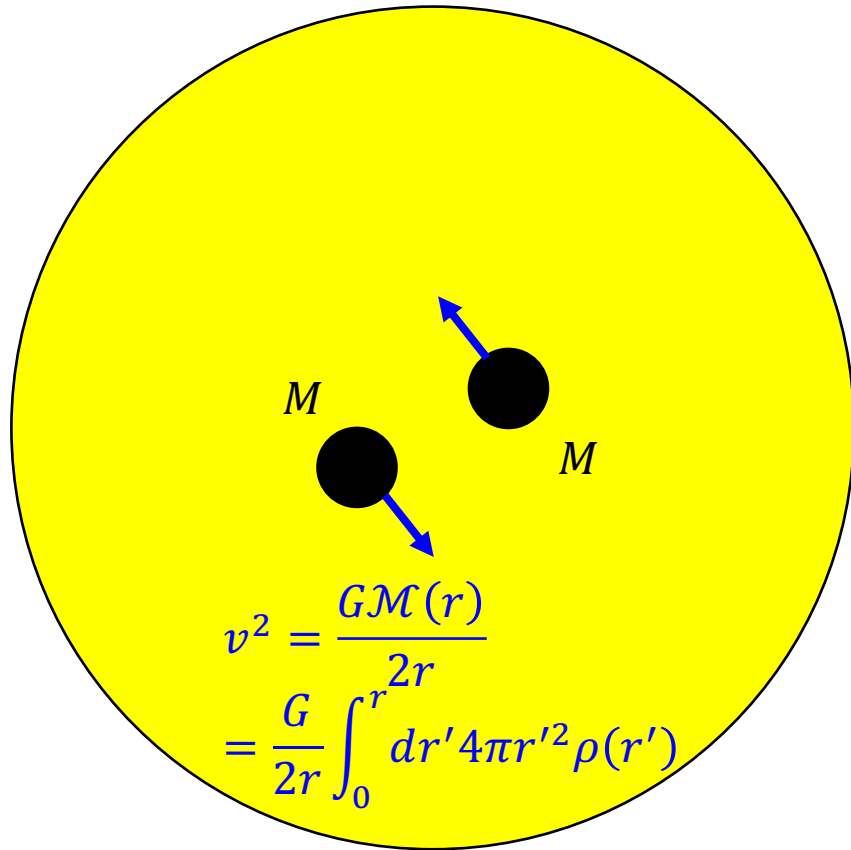
# Appendix 2 : Decay Timescale inside ULDM Halo

## 1. By Dynamical Friction



$$\begin{aligned}
 \tau &= \frac{L}{r \times |F_{\text{DF}}|} = \frac{Mvr}{r \times 4\pi\rho(r)(GM/v)^2 C(\tilde{r})} \\
 &= \frac{Mv^3}{4\pi\rho(r)(GM)^2 \times \frac{1}{3} \left(\frac{mvr}{\hbar}\right)^2} = \frac{3\hbar^2 \sqrt{\mathcal{M}(r)}}{8\pi m^2 \rho(r) M \sqrt{2G^3 r^5}}
 \end{aligned}$$

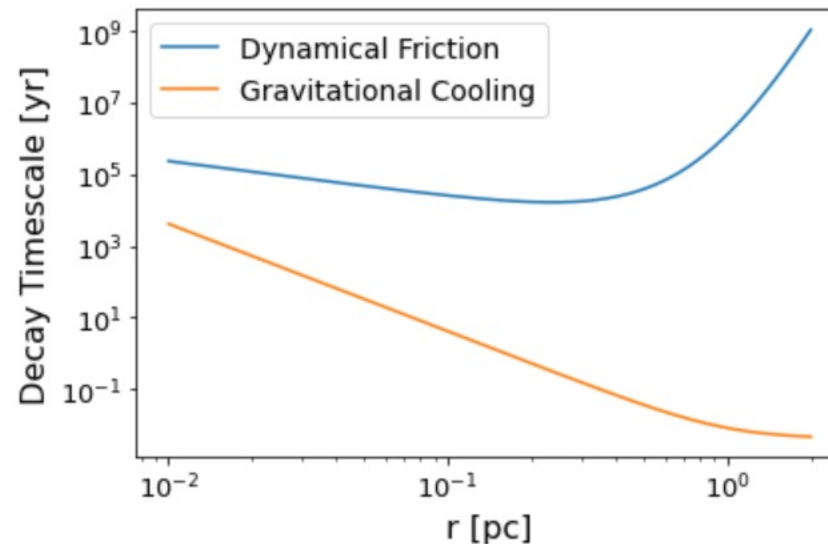
# Appendix 2 : Decay Timescale inside ULDM Halo



## 2. By Gravitational Cooling

$$D(t) = D_0 (1 + 3QD_0^3 t)^{-\frac{1}{3}} \Rightarrow \tau = \left( \frac{1}{D} \frac{dD}{dt} \right)^{-1} = \frac{1}{QD^3}$$

$$\text{Where } Q = \kappa \left( \frac{m}{10^{-21} \text{eV}/c^2} \right)^9 \frac{(M_s + 2\gamma_l M_{\text{bh}})^6}{M_{\text{bh}}} \left\{ \begin{array}{l} \kappa = \frac{29.58 \pm 1.39}{[\text{kyr} \cdot \text{pc}^3 \cdot (10^{10} M_\odot)^5]} \\ \gamma_l = 1.476 \pm 0.067 \cong 1.5 \end{array} \right.$$



$$\begin{aligned}
 m &= 10^{-21} \text{eV} c^2 \\
 M_s &= 10^{10} M_\odot \\
 M_{\text{bh}} &= 10^7 M_\odot
 \end{aligned}$$